Graphical Representations of Some Statistical Models

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- Introduction to critical phenomena
- High-temperature Graph
- Low-temperature Graph
- Fortuin-Kasteleyn Representation
- Baxter-Kelland-Wu Representation and
- Stochastic Lowner Evolution

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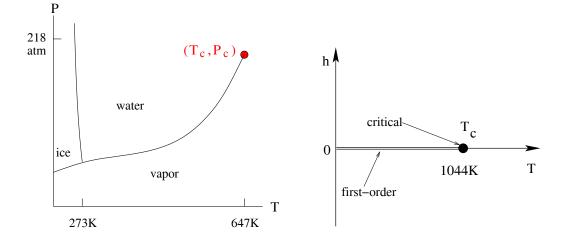
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Introduction





The liquid-gas critical point of H₂O: $T_c = 647K$, point of Fe: $h_c = 0$, $p_c = 218 \text{ atm.}$

The ferromagnetic $T_c = 1044K.$

- Statistical Models
 - Ising model on Graph (V, E).
 - 1, Configuration Space $\vec{s} \in \{+1, -1\}^{|V|}$.
 - 2, A priori measure $\frac{1}{2}(\delta_{s,+1} + \delta_{s,-1})$.
 - 3, Hamiltonian

$$\mathcal{H}/k_{\mathsf{B}}T = -K\sum_{\langle i,j\rangle} s_i s_j - H\sum_k s_k$$

- Partition function $Z = \sum e^{-\mathcal{H}/k_b T}$

or free energy $F = -\ln Z$

- Physical quantities as derivatives of F.

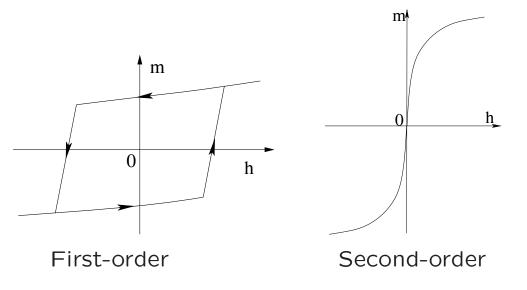
First derivative: ρ_{H_2O} , m

Second derivative: C, χ

 \star *n*th-order transition:

nth derivative of F is singular,

but (n-1)th derivative is analytic.

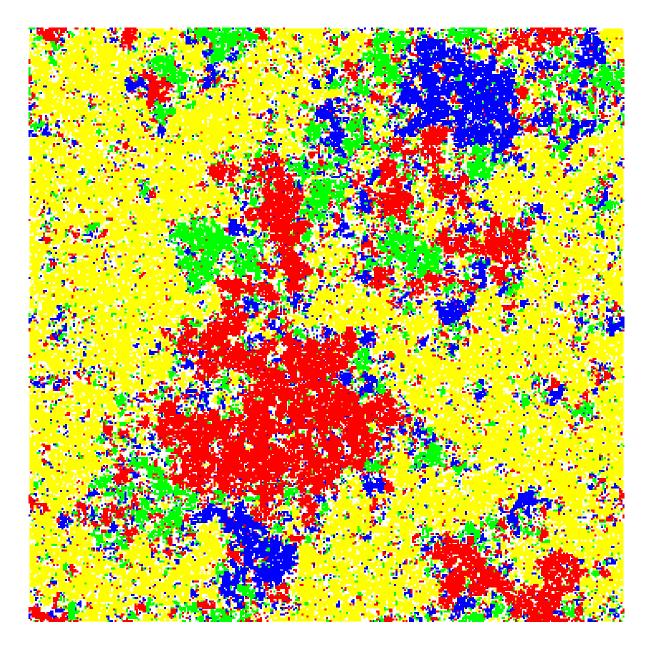


• Critical phenomena

Specific-heat : $C \propto |T - T_c|^{-\alpha}$ Susceptibility : $\chi \propto |T - T_c|^{-\gamma}$ Magnetization: $m \propto |T - T_c|^{\beta}$ At T_c , correlation: $g(r) \propto r^{-2X}$

***** Universality: critical exponents, α, γ, \cdots ,

take same values in different systems.



Four-state Potts at criticality; the Hamiltonian reads

$$\mathcal{H}/k_{\mathsf{B}}T = -J\sum_{\langle i,j\rangle}\delta_{\sigma_i,\sigma_j}$$

High-temperature expansion of the zero-field Ising model

Parition function:

$$Z = \sum_{\{\pm 1\}^{|V|}} e^{-\mathcal{H}}$$
$$= \sum_{\{\pm 1\}^{|V|}} \prod_{\langle i,j \rangle} e^{Ks_i s_j}$$
(1)

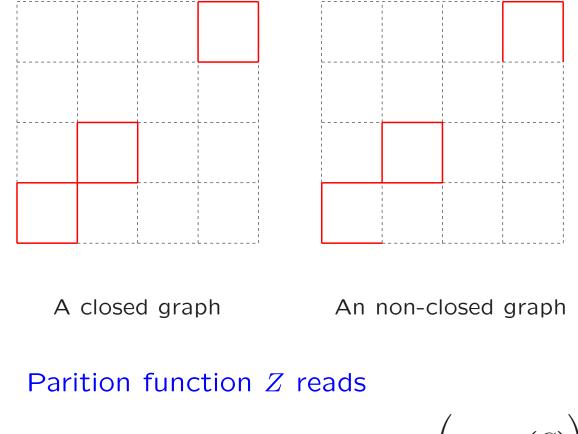
Identity

$$e^{Ks_is_j} = \cosh K + s_is_j \sinh K$$

= $\cosh K(1 + vs_is_j) \quad (v = \tanh K)$
(2)

Insert (2) into (1), and expand the product \prod . Then, label each term in the expansion by a graph Λ :

- If factor 1 is taken, do nothing;
- If vs_is_j , put a bond $\langle i,j \rangle$.



$$Z = (\cosh K)^{|E|} \sum_{G \in \mathcal{G}_0} v^{b(G)} \sum_{\{\pm 1\}^{|V|}} \left(\prod_i s_i^{m_i(G)} \right)$$

 $-m_i(G)$: number of bonds incident on vertex *i*.

*: Any graph with at least one *odd* number m_i contributes to Z by zero.

Z becomes

$$Z = 2^{|V|} (\cosh K)^{|E|} \sum_{G \in \mathcal{G}_0, \partial G = 0} v^{b(G)}.$$

★ High-T graph of $\langle s_i s_j \rangle$: Almost closed graphs with $\partial G = (i, j)$.

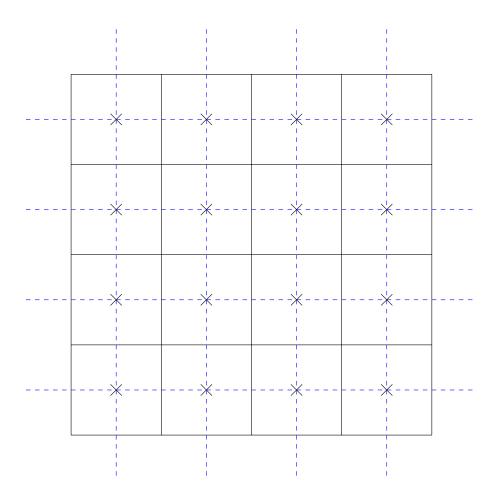
★ High-T graph of the Potts model: Flow polynomial

Application:

- 1D Ising model: $Z = 2^{L} (\cosh K)^{L}$ (free) $Z = 2^{L} (\cosh K)^{L} [1+(\tanh K)^{L}]$ (periodic)
- Exact solution of the 2D Ising model
- Worm Algorithm (has considerable applications in quantum systems).
- O(n) loop model and its cluster simulations.

Low-temperature expansion of the zero-field Ising model

Dual Lattice: $G \to G^*$

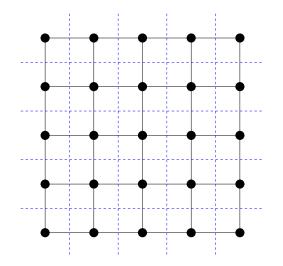


G: solid and black lines G^* : dashed and blue lines

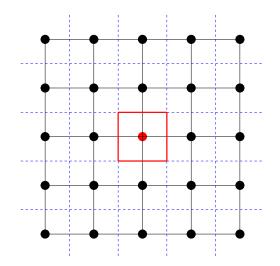
Partition function Z:

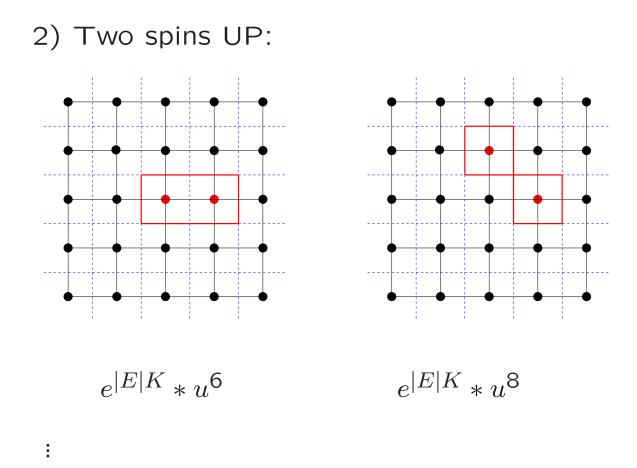
$$Z = \sum_{\{\pm 1\}^{|V|}} e^{-\mathcal{H}}$$

1) All spins DOWN: $e^{|E|K}$



2) One spin UP: $e^{|E|K} * e^{-8K} = e^{|E|K} * u^4$





Z becomes

$$Z = e^{|E|K} \sum_{G^* \in \mathcal{G}_0^*, \partial G^* = 0} u^{b(G^*)}.$$

Duality Relation:

 $Z(\mathcal{G}; v) = \text{constant} \times Z(\mathcal{G}^*; u)$ if

v = u. Namely, $\tanh K = e^{-2K^*}$

Criticality on Square Lattice: $2K_c = \ln(1 + \sqrt{2})$.

Fortuin-Kasteleyn Representation of the zero-field Ising model

Hamiltonian \mathcal{H}

$$\mathcal{H}/=-K\sum_{\langle i,j
angle}s_is_j=-2K\sum_{\langle i,j
angle}\delta_{s_i,s_j}+ ext{constant}$$

Identity:

$$e^{2K\delta_{s_i,s_j}} = 1 + \mu\delta_{s_i,s_j} \quad (\mu = e^{2K} - 1)$$

Partition function Z:

$$Z = \sum_{\{\pm 1\}^{|V|}} e^{-\mathcal{H}}$$
$$= \sum_{\{\pm 1\}^{|V|}} \prod_{\langle i,j \rangle} (1 + \mu \delta_{s_i,s_j})$$
(3)

Expand the product \prod , label each term by the graph: if the factor 1 is taken, do nothing; if $\mu \delta_{s_i,s_j}$ is taken, put a bond on $\langle i,j \rangle$.

 $\delta_{s_i,s_j} \Rightarrow \mathsf{Each}\ \mathsf{component}\ \mathsf{is}\ \mathsf{of}\ \mathsf{a}\ \mathsf{unique}\ \mathsf{color}$

Z becomes

$$Z = \sum_{G \in \mathcal{G}_0} \mu^{b(G)} 2^{k(G)}$$

-b: bond number, k: component number

Random-cluster model:

$$Z(q,\mu) = \sum_{G \in \mathcal{G}_0} \mu^{b(G)} q^{k(G)}$$

Some special cases:

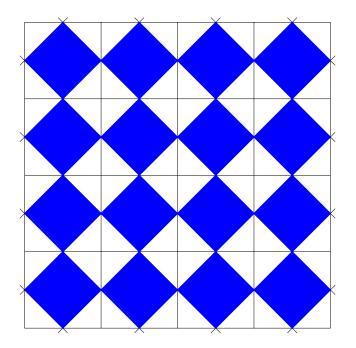
- $q \rightarrow 0$: spanning tree or spanning forests
- $\mu = -1$: Chromatic polynomial (such as The Four Color Prolem)

Monte Carlo Methods:

- Sweeny algorithm
- Swendsen-Wang-Chayes-Machta Algorithm

Baxter-Kelland-Wu (BKW) mapping of the random-cluster model

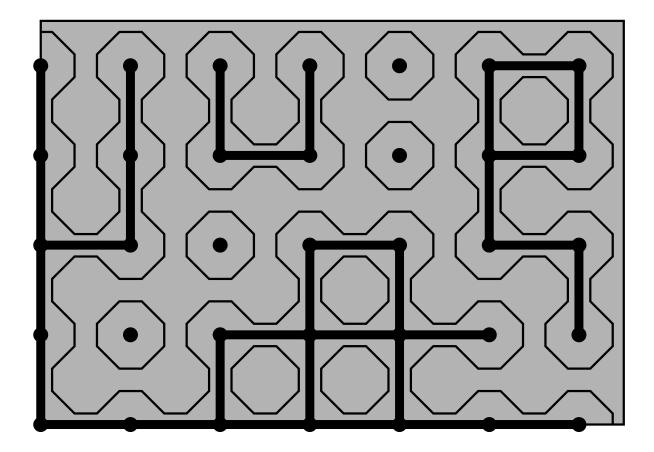
Medial Graph : $G \to \tilde{G}$



 $\tilde{G}-\mbox{Water:}$ blue area; Island: white

Operation:





Random-Cluster (RC) Model:

$$Z(q,\mu) = \sum_{G \in \mathcal{G}_0} \mu^{b(G)} (\sqrt{q})^{c(G)}$$

-c(G): loop number

Application: Map the RC model onto the 8-vertex model.

Recent Development: Stochastic Lowner Evolution (SLE), by Oded Schramm (1999).

In the scaling limit, critical systems are not only *scale invariant*, but also conformally invariant.

In the scaling limit, BKW curves with appropriate boundary conditions are described by SEL_{κ} in Z^+ :

$$\frac{\partial g_t(z)}{\partial t} = \frac{2}{g_t(z) - \sqrt{\kappa}B_t} , g_0(z) = z ,$$

— $B_t, t \in [0, \infty)$: Brownian motion.

