

关于Ampère环路定理的讨论

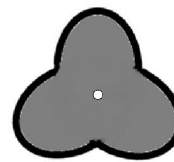


This is how scientists see the world.
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A simple example



平面“三叶草”回路



$$\oint_C \mathbf{B} \cdot d\mathbf{l} = \mu_0 \iint_S \mathbf{J} \cdot d\mathbf{S} = \mu_0 I \quad ?$$



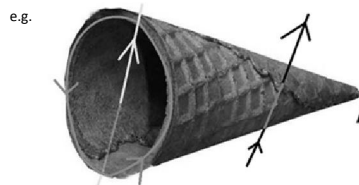
Parametric Equation:

$$\begin{cases} x = \sin t + 2\sin 2t \\ y = -\cos t + 2\cos 2t \end{cases} \quad t: 2\pi \rightarrow 0$$

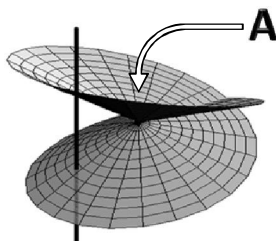
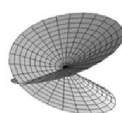
$$\mathbf{B} = \frac{\mu_0 I \Phi}{2\pi r}, \quad \mathbf{r} = (x, y) \Phi \text{ 为环向单位矢量 } \Phi = \frac{(-y, x)}{r}$$

$$\begin{aligned} \oint_C \mathbf{B} \cdot d\mathbf{l} &= \frac{\mu_0 I}{2\pi} \oint_C \frac{(-y, x)}{r^2} \cdot d\mathbf{l} \\ &= \frac{\mu_0 I}{2\pi} \int_C \frac{-y dx + x dy}{x^2 + y^2} \\ &= \frac{\mu_0 I}{2\pi} \int_{2\pi}^0 \frac{(\cos t - 2\cos 2t)(\cos t + 4\cos 2t) dt + (\sin t + 2\sin 2t)(\sin t - 4\sin 2t) dt}{(\sin t + 2\sin 2t)^2 + (-\cos t + 2\cos 2t)^2} \\ &= \frac{\mu_0 I}{2\pi} \int_{2\pi}^0 \frac{2\cos(t)\cos(2t) - 2\sin(t)\sin(2t) - 7}{4\sin(t)\sin(2t) - 4\cos(t)\cos(2t) + 5} dt \\ &= \frac{\mu_0 I}{2\pi} \cdot 4\pi \\ &= 2\mu_0 I \end{aligned}$$

Definition: 任取空间中一点A，曲线上一动点绕曲线运动一周与A连线所扫过的面，定义为该曲线围成的曲面。

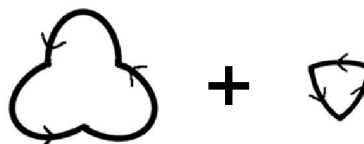


Remark: 对静磁学A点确实是任意的，正如Stokes' theorem中的曲面的任意性，而任意的好处使得计算可以避免奇点。



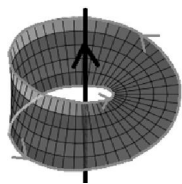
3D
↓
Z → 0 Degenerate
2D

小技巧

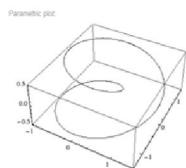


并不具有一般性!

II. The Boundary of Möbius Strip



①如图，电流I穿过Möbius带中间的空隙，求绕Möbius带边界的磁环量。



Parametric Equation (midcircle of radius 1):

$$x = (1 + \frac{1}{2} \cos \frac{u}{2}) \cos u$$

$$y = (1 + \frac{1}{2} \cos \frac{u}{2}) \sin u$$

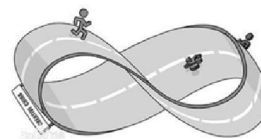
$$z = \frac{1}{2} \sin \frac{u}{2}$$

$$0 \leq u < 4\pi$$

小技巧已经行不通了!

证1: Möbius曲面不定向, 所以无论如何拆成一块块小环路, 总会出现添加的相邻路径无法抵消的情况, 否则由每一小块回路按右手法则确定的方向一致导出Möbius曲面定向, 矛盾。

证2: 如果拆分法可行, 对每一小回路运用Stokes' Theorem, 由Möbius曲面上各点旋度为0 ($\nabla \times \frac{\mu_0 I}{2\pi r} \phi = 0$)可得磁环量为0



老老实实计算:

$$\mathbf{B} = \frac{\mu_0 I \phi}{2\pi r}, r=(x,y,0) \quad \Phi \text{为环向单位矢量} \Phi = \frac{(-y, x, 0)}{r}$$

$$\oint_C \mathbf{B} \cdot d\mathbf{l} = \frac{\mu_0 I}{2\pi} \oint_C \frac{(-y, x, 0)}{r^2} \cdot d\mathbf{l}$$

$$= \frac{\mu_0 I}{2\pi} \int_C -y dx + x dy$$

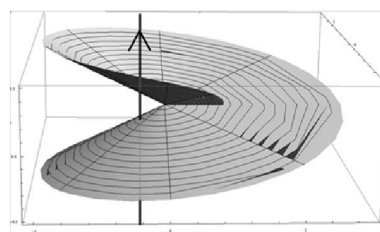
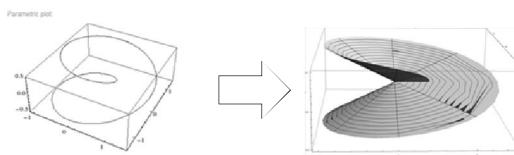
$$= \frac{\mu_0 I}{2\pi} \int_0^{4\pi} 1 \cdot du$$

$$= \frac{\mu_0 I}{2\pi} \cdot 4\pi$$

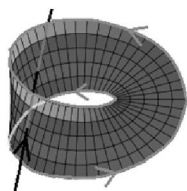
$$= 2\mu_0 I$$

WHY?

Möbius 带边界 围成的曲面 \Rightarrow Möbius 曲面



Also



②如图电流I穿过Möbius带, 求绕Möbius带边界的磁环量。

导线方程:

$$y = -\frac{1}{4}, z = 0$$

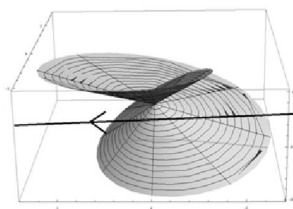
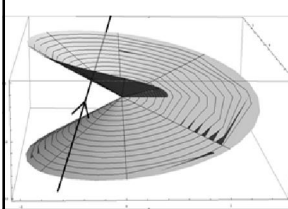
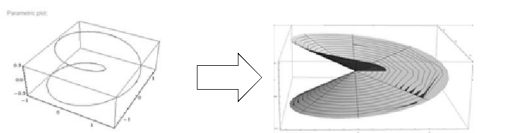
$$\mathbf{B} = \frac{\mu_0 I \phi}{2\pi r}, r=(0, y + \frac{1}{4}, z) \quad \Phi \text{为环向单位矢量} \Phi = \frac{(0, z, -y - \frac{1}{4})}{r}$$

$$\oint_C \mathbf{B} \cdot d\mathbf{l} = \frac{\mu_0 I}{2\pi} \oint_C \frac{(0, z, -y - \frac{1}{4})}{r^2} \cdot d\mathbf{l}$$

$$= \frac{\mu_0 I}{2\pi} \int_C z dy - (y + \frac{1}{4}) dz$$

$$= \frac{\mu_0 I}{2\pi} \int_0^{4\pi} \frac{z \sin(\frac{u}{2}) - ((1 + \frac{1}{2} \cos \frac{u}{2}) \sin u + \frac{1}{4}) \cos(\frac{u}{2})}{(\frac{1}{2} \sin \frac{u}{2})^2 + ((1 + \frac{1}{2} \cos \frac{u}{2}) \sin u + \frac{1}{4})^2} du$$

$$= 0$$



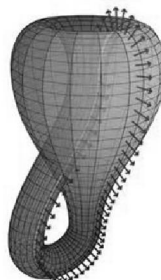
根本原因：不定向！

Non-orientable!

In the very definition of surface integral, integral can only be calculated on orientable surfaces.

To deal with such non-orientable or degenerated surfaces, our method is to redefine the exact surface of the circle, which is absolutely orientable.

Orientation also matters in Gauss's law



Klein Bottle

$$\oiint_S \mathbf{E} \cdot d\mathbf{S} = \iiint \nabla \cdot \mathbf{E} dV = 0$$

Trivial

谢谢观看

Xie Xie