

Atmospheric CO₂ capture and photofixation to near-unity CO by Ti³⁺-V_o-Ti³⁺ sites confined in TiO₂ ultrathin layers

Liang Liang^{1†}, Peiquan Ling^{1†}, Yuhuan Li^{1†}, Li Li¹, Jiandang Liu¹, Qiquan Luo^{2*},
Hongjun Zhang^{1*}, Qian Xu¹, Yang Pan¹, Junfa Zhu¹, Bangjiao Ye¹ & Yongfu Sun^{1*}

¹Hefei National Laboratory for Physical Science at Microscale, National Synchrotron Radiation Laboratory, University of Science and Technology of China, Hefei 230026, China;

²Institute of Physical Science and Information Technology, Anhui University, Hefei 230601, China

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To realize efficient atmospheric CO₂ chemisorption and activation, abundant Ti³⁺ sites and oxygen vacancies in TiO₂ ultrathin layers were designed. Positron annihilation lifetime spectroscopy and theoretical calculations first unveil each oxygen vacancy is associated with the formation of two Ti³⁺ sites, giving a Ti³⁺-V_o-Ti³⁺ configuration. The Ti³⁺-V_o-Ti³⁺ sites could bond with CO₂ molecules to form a stable configuration, which converted the endoergic chemisorption step to an exoergic process, verified by *in-situ* Fourier-transform infrared spectra and theoretical calculations. Also, the adjacent Ti³⁺ sites not only favor CO₂ activation into COOH* *via* forming a stable Ti³⁺-C-O-Ti³⁺ configuration, but also facilitate the rate-limiting COOH* scission to CO* by reducing the energy barrier from 0.75 to 0.45 eV. Thus, the Ti³⁺-V_o-TiO₂ ultrathinlayers could directly capture and photofix atmospheric CO₂ into near-unity CO, with the corresponding CO₂-to-CO conversion ratio of ca. 20.2%.

TiO₂ ultrathin layers, CO₂ activation, CO₂ photofixation, positron annihilation lifetime spectroscopy

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Recently, as monitored by the CO₂-tracking satellites, atmospheric concentrations of CO₂ reach an all-time high of 33 gigatonnes in 2019, which inevitably causes environmental deterioration and climate change [1]. Through mimicking the natural photosynthesis of plants, artificial photocatalysts have been widely fabricated for capturing and converting CO₂ into fuels and value-added products [2]. Although the artificial photosynthetic process may contribute to lower the levels of CO₂ in atmosphere, most of the previous studies require high purity (typically ≥99.99%) of CO₂, which is costly and arduous to be separated and concentrated from ambient air levels (~400 ppm) as well as compressed to typically 11–14 MPa for transportation and storage [3]. In

other words, only very limited success has been achieved in low concentration of CO₂ photoreduction, especially under atmospheric-like conditions (~0.04%) [4]. This can primarily be ascribed to thermodynamic stability and kinetic inertness of CO₂ molecule, since the dissociation energy of its C=O bond is as high as ~750 kJ mol⁻¹. More importantly, many of the reported photocatalysts still suffer from low selectivity for the desirable products [5], mainly because of the various possible CO₂ reduction pathways toward different products. To overcome these limitations, major considerations should be prioritised to efficiently capture CO₂ molecules and selectively convert them into the single product.

With regard to the initial CO₂ adsorption process, increasing the number of low-coordinated sites on the surface of catalysts could promote the ability of CO₂ chemisorption, since the surface sites with many dangling bonds can

†These authors contributed equally to this work.

*Corresponding authors (email: yfsun@ustc.edu.cn; hjzhang8@ustc.edu.cn; qluo@ahu.edu.cn)

strongly bind with the CO₂ molecules [6]. Concerning the following CO₂ activation procedure, creating abundant electron-donating sites is able to transform the linear and chemically inert CO₂ molecule to a bent structure, benefiting from the electron donation to the electron-poor carbon atoms [7]. Considering that the product selectivity is closely related to the intermediate configurations on the surface of catalysts [8], it is highly necessary to modulate the binding nature of key reaction intermediates and hence provide the possibility to tailor product selectivity. Taken together, designing a photocatalyst featuring atomic-level thickness and abundant surface defects is highly desirable, where the atomic thickness can maximize the number of surface sites for CO₂ adsorption, while the defect-induced electron localization centers could act as active sites for activating inert CO₂ molecules.

Herein, taking the most widely used photocatalyst of TiO₂ as an example, defective anatase TiO₂ ultrathin layers were first successfully synthesized using the mixed-phase TiO₂ ultrathin layers with a scale of ~1.50 g/batch as the precursor (Figure S1, Supporting Information online). X-ray diffraction (XRD) pattern of the products in Figure 1(a) could be readily indexed to anatase TiO₂, corresponding to JCPDS No. 21-1272. All the peaks in the micro-Raman spectrum (Figure S2) could further verify that the TiO₂ products were

in pure anatase phase, where no peaks of any other phases or carbon impurities in the range of 100–1700 cm⁻¹ were detected. Moreover, transmission electron microscopy (TEM) image in Figure 1(b) exhibited their freestanding and sheet-like morphology with a lateral size of ca. 200 nm. High-resolution TEM image in Figure S3(a) depicted that the interplanar distances of 2.38 and 3.52 Å corresponded to the (004) and (101) lattice plane distances of the anatase TiO₂. Meanwhile, their dihedral angle of 69° was also fairly consistent with the calculated angle between (004) and (101) planes, implying the [010] orientation of the synthesized TiO₂ ultrathin layers. Atomic force microscopy (AFM) and the corresponding height profiles in Figure S3(b, c) clearly manifested that the average thickness of the as-obtained ultrathin layers was ca. 1.42 nm, which agreed well with the thickness of 7-atom thick TiO₂ slab along the [010] direction. More importantly, as revealed by the X-ray photoelectron spectra (XPS) in Figure 1(c), the Ti 2p spectra could be deconvoluted into two doublets corresponding to the Ti⁴⁺ (2p_{3/2}=458.8 eV) and Ti³⁺ (2p_{3/2}=457.4 eV), whereas the O 1s core-level spectra could be fitted by three peaks corresponding to O_{lattice} (530.1 eV), O_{Ti³⁺} (531.1 eV) and oxygen vacancy (531.4 eV), indicating the formation of Ti³⁺ and oxygen vacancies in the synthesized TiO₂ ultrathin layers [9]. By comparison, the pristine TiO₂ ultrathin layers were also

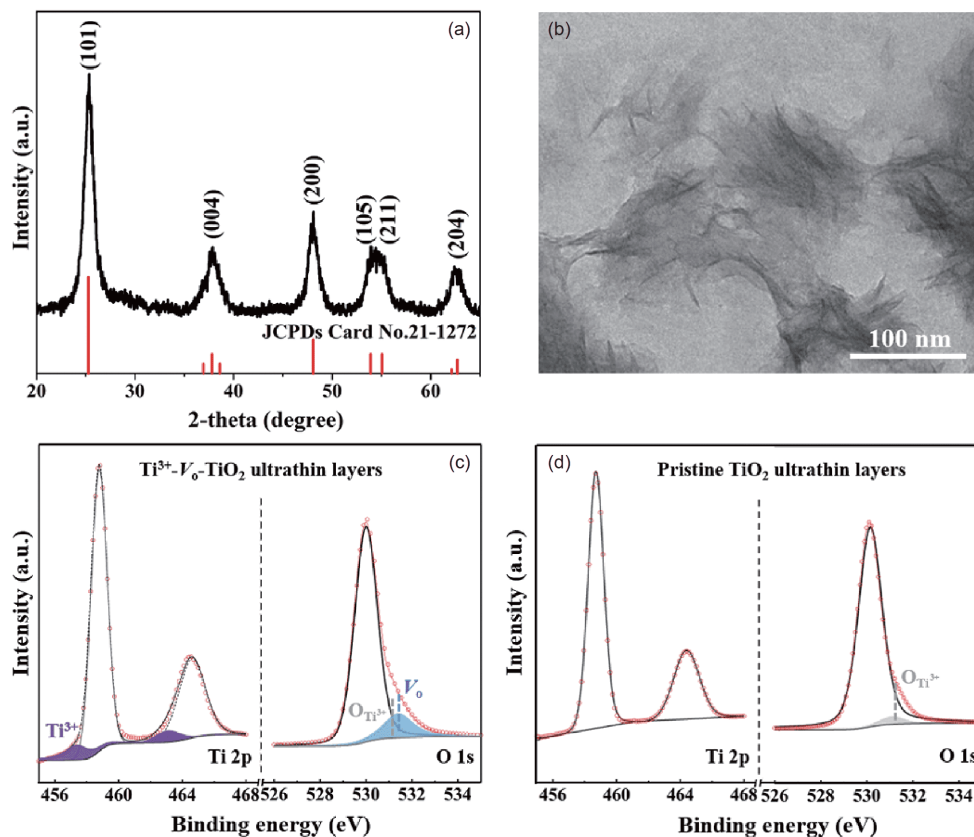


Figure 1 (a) XRD pattern, (b) TEM image for the Ti³⁺-V_o-TiO₂ ultrathin layers; (c) Ti 2p and O 1s XPS spectra for the Ti³⁺-V_o-TiO₂ ultrathin layers; (d) Ti 2p and O 1s XPS spectra for the pristine TiO₂ ultrathin layers (color online).

fabricated by tuning the annealing atmosphere (Figure S4), in which no obvious signals for Ti^{3+} and oxygen vacancies were detected by the corresponding XPS spectra in Figure 1 (d), confirming the negligible amount of defect structure within the pristine TiO_2 ultrathin layers.

To further identify the Ti^{3+} and oxygen vacancies in the synthesized defective TiO_2 ultrathin layers, density functional theory (DFT) calculations and positron annihilation lifetime (PAL) spectroscopy were adopted [10,11]. Initially, DFT calculations in Figure S5 illustrated that the bridging oxygen vacancy was easier to be formed within the defective TiO_2 layer slab as compared to other types of oxygen vacancies. In addition, Bader charges of the two nearest-neighbor Ti atoms adjacent to the bridging oxygen vacancy presented 1.79 and 1.75 electrons (Figure 2(a, b)), respectively, which implied their +3 oxidation state by comparing with that of the Ti^{3+} (1.79e) in Ti_2O_3 [12]. Furthermore, as illustrated by the PAL spectra in Figure 2(c, d), Figure S6 and Table 1, the defective TiO_2 ultrathin layers and the pristine TiO_2 ultrathin layers showed three lifetime components. As compared to the calculated positron lifetime values (186.8 ps for defect-free, 192.4 ps for $\text{Ti}^{3+}\text{-V}_\text{o}\text{-Ti}^{3+}$, 195.4 ps for $\text{V}_\text{o}\text{-o}$, 197.5 ps for V_Ti , 209.5 ps for $\text{V}_\text{Ti-O}$, 203.1 ps for $\text{V}_\text{Ti-Ti}$) and Figure S7, the experimental shortest positron lifetime (τ_1) for both the samples corresponded to positron annihilation at the $\text{Ti}^{3+}\text{-V}_\text{o}\text{-Ti}^{3+}$ configuration sites and defect-free region, respectively [11]. The other two longer components (τ_2 and τ_3) for both the samples could be attributed to the large defect clusters and the interface present in the samples, respectively [13]. As an outcome, all the above theoretical and experimental results undoubtedly demonstrated the formation of $\text{Ti}^{3+}\text{-V}_\text{o}\text{-Ti}^{3+}$ defective structure in the synthetic TiO_2 ultrathin layers ($\text{Ti}^{3+}\text{-V}_\text{o}\text{-TiO}_2$ ultrathin layers).

To disclose the effects of $\text{Ti}^{3+}\text{-V}_\text{o}\text{-Ti}^{3+}$ defective structure on the electronic band structure, synchrotron-radiation photoelectron spectra (SRPES) and ultraviolet-visible spectroscopy (UV-Vis) diffuse reflectance spectra were conducted on the $\text{Ti}^{3+}\text{-V}_\text{o}\text{-TiO}_2$ ultrathin layers and the pristine TiO_2 ultrathin layers [14]. As shown in Figure S8(a), the $\text{Ti}^{3+}\text{-V}_\text{o}\text{-TiO}_2$ ultrathin layers and the pristine TiO_2 ultrathin layers exhibited the band gaps of 3.30 and 3.35 eV, respectively. As clarified by SRPES spectra in Figure S8(b), the $\text{Ti}^{3+}\text{-V}_\text{o}\text{-TiO}_2$ ultrathin layers and the pristine TiO_2 ultrathin layers possessed the work function of 4.20 and 4.25 eV, the valence band maximums of 2.97 and 3.1 V vs. normal hydrogen electrode (NHE), and the conduction band minimums of -0.33 and -0.25 V vs. NHE, respectively. Based on the

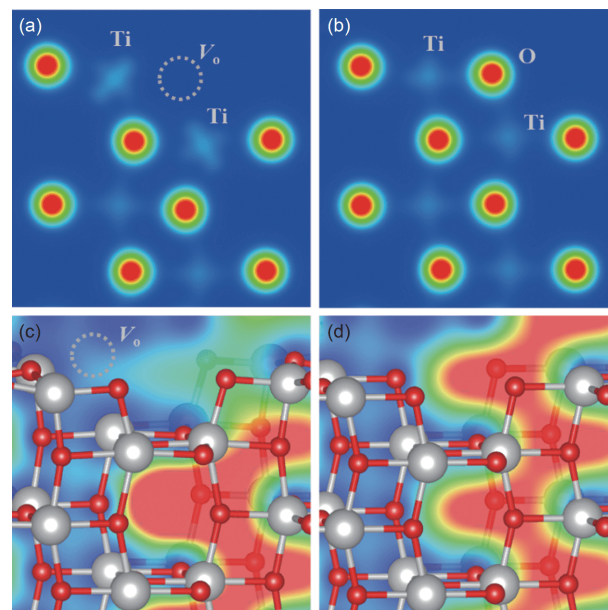


Figure 2 Distribution of charge density for (a) the $\text{Ti}^{3+}\text{-V}_\text{o}\text{-TiO}_2$ ultrathin layer slab and (b) the pristine TiO_2 ultrathin layer slab, plotted from 0 (blue) to $0.035 \text{ e Bohr}^{-3}$ (red). Distribution of positron density for (c) the $\text{Ti}^{3+}\text{-V}_\text{o}\text{-TiO}_2$ ultrathin layer slab and (d) the pristine TiO_2 ultrathin layer slab. Bader charges for the Ti sites adjacent to the oxygen vacancy of $\text{Ti}^{3+}\text{-V}_\text{o}\text{-TiO}_2$ ultrathin layer slab are 1.79e and 1.75e, respectively, showing a +3 oxidation state. The grey and red balls in (c, d) represent the Ti and O atoms, respectively (color online).

above results, the band structures for the $\text{Ti}^{3+}\text{-V}_\text{o}\text{-TiO}_2$ ultrathin layers and the pristine TiO_2 ultrathin layers theoretically matched with the CO_2 reduction and H_2O oxidation potentials simultaneously (Table S1, Supporting Information online), providing the prerequisite to participate in the following reactions.

As is well-known, the rising atmospheric levels of CO_2 raise concerns about the ensuing effects on the global climate. In this respect, conversion of atmospheric CO_2 to hydrocarbons could be expected as a promising way to reduce CO_2 emissions and gain value-added products simultaneously [15]. Considering the unique defective structure and well-matched band structure for TiO_2 ultrathin layers, their CO_2 photoreduction experiments were carried out in the fresh air, simulated air and pure CO_2 , respectively. Expectedly, in the fresh air, the $\text{Ti}^{3+}\text{-V}_\text{o}\text{-TiO}_2$ ultrathin layers possessed a CO evolution rate of $23.1 \mu\text{mol g}^{-1} \text{ h}^{-1}$ (Figure 3 (a)), which was roughly 2.8 times higher than that of the pristine TiO_2 ultrathin layers; meanwhile, only negligible amount of H_2 was detected and no any other gas or liquid carbonous product was detected by gas chromatography or

Table 1 Analysis results of experimental PAL spectra for the $\text{Ti}^{3+}\text{-V}_\text{o}\text{-TiO}_2$ ultrathin layers and the pristine TiO_2 ultrathin layers, respectively

	τ_1 (ps)	τ_2 (ps)	τ_3 (ns)	I_1 (%)	I_2 (%)	I_3 (%)
$\text{Ti}^{3+}\text{-V}_\text{o}\text{-TiO}_2$ ultrathin layers	192.4 ± 2.2	350.7 ± 0.3	2.150 ± 0.030	10.8 ± 0.2	86.1 ± 0.2	3.1 ± 0.1
Pristine TiO_2 ultrathin layers	188.7 ± 2.9	399.2 ± 0.4	1.616 ± 0.016	6.9 ± 0.1	90.0 ± 0.1	3.1 ± 0.1

^1H nuclear magnetic resonance (NMR) during the photocatalysis in the fresh air, thus giving a CO selectivity of near 100%. More importantly, the $\text{Ti}^{3+}\text{-V}_\text{o}\text{-TiO}_2$ ultrathin layers showed the CO_2 -to-CO conversion ratio of ca. 20.2% after 9 h photocatalysis in the fresh air (Figure 3(b)), which was obviously higher than the 0.01% conversion ratio in pure CO_2 , giving possibilities of the direct implementation of CO_2 in the fresh air. $^{13}\text{CO}_2$ isotope-labelling experiments combined with the synchrotron-radiation vacuum ultraviolet photoionization mass spectrometry (SVUV-PIMS) (Figure 3(c)) further identified the product of ^{13}CO was indeed originated from the reduction of isotope-labeled $^{13}\text{CO}_2$. *In-situ* thermographic photographs in Figure S9 depicted that the $\text{Ti}^{3+}\text{-V}_\text{o}\text{-TiO}_2$ ultrathin layers could maintain the temperature at $\sim 26^\circ\text{C}$ during a catalytic cycle, which demonstrated they could achieve CO_2 photoreduction at room temperature in the fresh air [16]. Furthermore, the $\text{Ti}^{3+}\text{-V}_\text{o}\text{-TiO}_2$ ultrathin layers did not show any obvious attenuation in CO_2 photoreduction activities in the fresh air even after the fourth cycle of repetition test up to 36 h, while the post-catalysis characterizations showed their superior photostability (Figure 3(d) and Figure S10). It is worth noting that these results tested in the fresh air were consistent with that in the simulated air (21.96% O_2 + 78% N_2 + 0.04% CO_2), further verifying the $\text{Ti}^{3+}\text{-V}_\text{o}\text{-TiO}_2$ ultrathin layers could capture atmospheric CO_2 and convert it into valuable products. With regard to the other half reaction, the O_2 evolution rate was

measured to be ca. $10.1\ \mu\text{mol g}^{-1}\text{ h}^{-1}$ after subtraction of the experimental background concentration of O_2 in the air for the $\text{Ti}^{3+}\text{-V}_\text{o}\text{-TiO}_2$ ultrathin layers, which was consistent with the expected stoichiometric ratio (2:1) of CO/O_2 ratio for CO_2 reduction. Taken together, the above results verified that the $\text{Ti}^{3+}\text{-V}_\text{o}\text{-TiO}_2$ ultrathin layers could achieve 20.2% CO_2 -to-CO conversion and $\sim 100\%$ CO selectivity in atmospheric CO_2 photoreduction, hence providing a feasible pathway for practical solar fuel applications.

To better understand why the $\text{Ti}^{3+}\text{-V}_\text{o}\text{-Ti}^{3+}$ defective sites contributed to improved conversion and selectivity for atmospheric CO_2 reduction, the capture and activation ability of CO_2 for the $\text{Ti}^{3+}\text{-V}_\text{o}\text{-TiO}_2$ ultrathin layers was compared with that of the pristine TiO_2 ultrathin layers. Initially, as displayed by the CO_2 adsorption isotherms in Figure S11, the $\text{Ti}^{3+}\text{-V}_\text{o}\text{-TiO}_2$ ultrathin layers revealed an increased CO_2 adsorption capacity as compared to the pristine TiO_2 ultrathin layers, indicating the stronger adsorption ability on the former sample. To further study the CO_2 adsorption configuration on the surface of $\text{Ti}^{3+}\text{-V}_\text{o}\text{-TiO}_2$ ultrathin layers, *in-situ* Fourier transformed infrared spectroscopy (FTIR) was collected under gas flow of the simulated air and at the presence of water. The peak at $1507\ \text{cm}^{-1}$ in Figure 4(a) and Figure S12(a) revealed the presence of monodentate carbonate species, while the peak centered at $1672\ \text{cm}^{-1}$ could be assigned to the bidentate carbonate species, confirming the carbonate species of CO_2 adsorption configuration on the

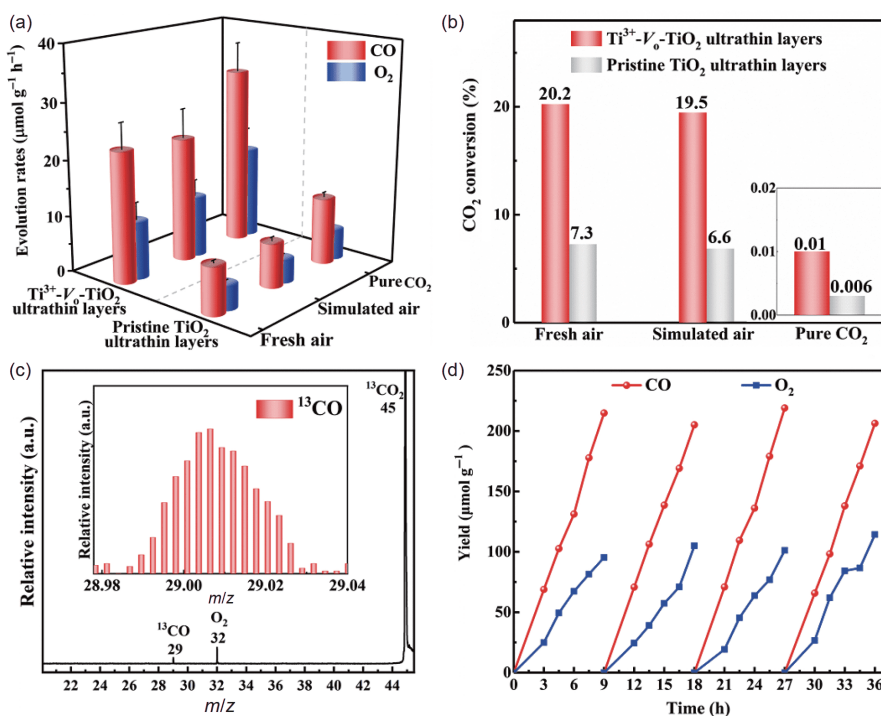


Figure 3 (a) CO_2 photoreduction experiments carried out in the fresh air, simulated air and pure CO_2 , respectively, and (b) the corresponding CO_2 -to-CO conversion ratio over the $\text{Ti}^{3+}\text{-V}_\text{o}\text{-TiO}_2$ ultrathin layers and the pristine TiO_2 ultrathin layers. (c) SVUV-PIMS spectrum of the products for the $\text{Ti}^{3+}\text{-V}_\text{o}\text{-TiO}_2$ ultrathin layers after $^{13}\text{CO}_2$ photoreduction at $h\nu=14.50\ \text{eV}$. Inset gives the corresponding signal of $m/z=29$ (^{13}CO). (d) Cycling curves of CO and O_2 evolution rates over the $\text{Ti}^{3+}\text{-V}_\text{o}\text{-TiO}_2$ ultrathin layers. The error bars in (a) represent the standard deviations of three independent measurements of the same sample (color online).

surface [17]. To gain insights into the geometric effect and electronic effect of CO₂ adsorption over the Ti³⁺-V_o-TiO₂ ultrathin layers, the DFT simulations and charge density difference were adopted. As shown in Figure S13, the angle of Ti–O–Ti in the same row as the V_o sites increased from 99.98° to 101.50°, indicating that the geometric effect of V_o sites led to surface relaxation. Once CO₂ is adsorbed, the structure was counteracted with the decrease angle of Ti–O–Ti of 101.12° by achieving a bent CO₂ configuration. Notably, there was charge accumulation in the bonding region between CO₂ and the Ti³⁺-V_o-Ti³⁺ defective sites (Figure S14), which suggested the strong interaction between CO₂ molecules and the defective sites resulted in originally linear CO₂ molecule's bent configuration with an O–C–O angle (θ in Figure S14) of 135.91° and two elongated C–O bonds of 1.24 and 1.27 Å, respectively, thus forming a relative stable configuration (Figure S15). In addition, the adsorption energy was estimated to be –0.10 eV, further indicating the strong bonding between CO₂ and the Ti³⁺-V_o-Ti³⁺ defective sites. As comparison, CO₂ adsorption energy on the pristine TiO₂ ultrathin layer slab was calculated to be slightly endothermic with an adsorption energy of +0.02 eV, indicating

the pristine surface was relatively difficult to bond with CO₂ due to their intrinsic chemical inertness. The above-mentioned results showed that the exposed Ti³⁺-V_o-Ti³⁺ defective sites with the presence of excess electrons helped to strengthen the bonding between CO₂ and Ti³⁺-V_o-Ti³⁺ sites, and thus achieved improved CO₂ adsorption ability. Moreover, the Ti³⁺-V_o-TiO₂ ultrathin layers exhibited the decreased intensity of the photoluminescence emission spectra (Figure S16) and increased surface photovoltage (SPV) response intensity (Figure S17) relative to the pristine TiO₂ ultrathin layers, suggesting the Ti³⁺-V_o-Ti³⁺ defective sites were beneficial for photogenerated electron-hole separation, and hence able to afford more electrons to activate CO₂ molecules [18].

To further investigate the CO₂ activation process, *in-situ* FTIR spectra were performed to probe the reaction intermediates over the Ti³⁺-V_o-TiO₂ ultrathin layers. As shown in Figure 4(b) and Figure S12(b), a distinct peak appeared at ca. 1622 cm⁻¹ could be assigned to hydrocarboxyl (COOH*) intermediate, suggesting the CO₂ could be reduced into CO *via* the COOH* intermediates [19]. Based on the above experimental results, the DFT calculations (Figure 4(c), Figure

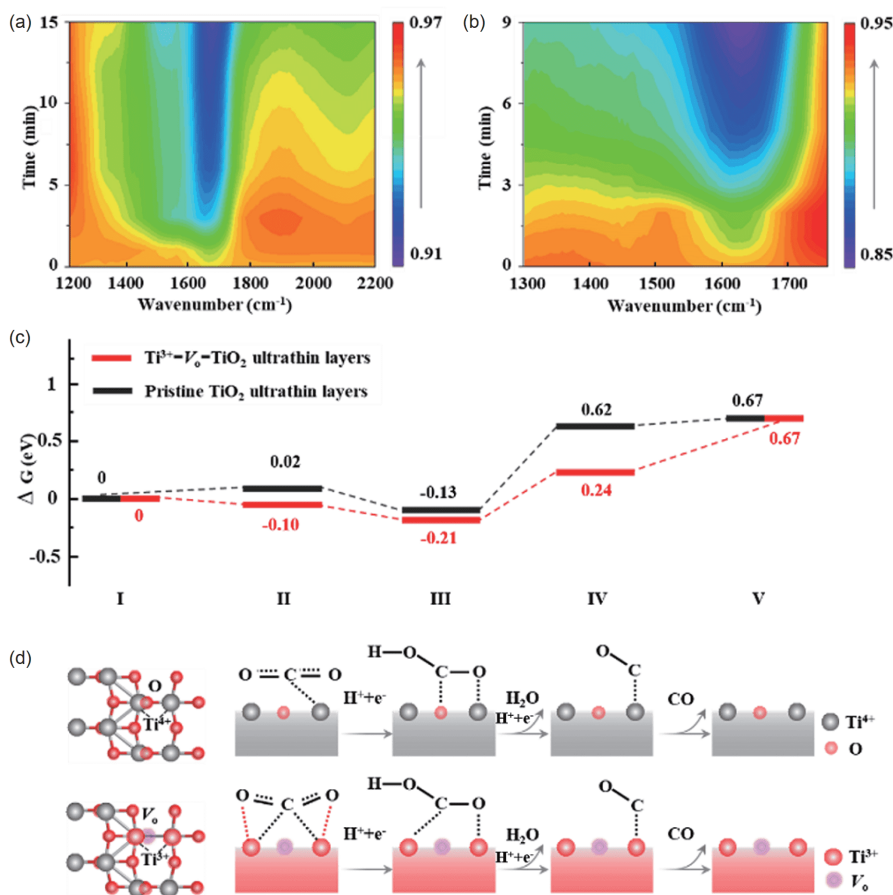


Figure 4 (a) 2D contour map of *in-situ* FTIR spectra for the CO₂ adsorption over the Ti³⁺-V_o-TiO₂ ultrathin layers; (b) 2D contour map of *in-situ* FTIR spectra for the reaction intermediates during the CO₂ photoreduction process over the Ti³⁺-V_o-TiO₂ ultrathin layers; (c) free energy diagrams of CO₂ photoreduction and (d) schemes for the corresponding key elementary steps (I–V) over the Ti³⁺-V_o-TiO₂ ultrathin layer slab and the pristine TiO₂ ultrathin layer slab, respectively (color online).

S18 and Table S2) were performed to simulate the process of CO₂ reduction. In the step of the protonation of the adsorbed CO₂ to form an adsorbed COOH* intermediate (Figure 4(d)), two Ti³⁺ sites in the Ti³⁺-V_o-Ti³⁺ defective sites could bond with COOH* to form a bridge configuration of Ti³⁺-C-O-Ti³⁺, resulting in a lowered energy barrier for COOH* formation. Note that the subsequent step of COOH* scission to CO* was regarded as the rate-limiting step, which showed an energy barrier of 0.45 eV, obviously lower than the 0.75 eV for the pristine TiO₂ ultrathin layer slab. Besides, the Ti³⁺ sites on the Ti³⁺-V_o-TiO₂ ultrathin layer slab helped to stabilize the CO* intermediate with a stronger CO binding energy of -0.43 eV as compared to that of the pristine TiO₂ layer slab (-0.05 eV), suggesting that the formation of Ti³⁺-CO provided an additional driving force for the scission step of COOH* to CO*. Accordingly, the Ti³⁺-V_o-Ti³⁺ defective sites in the TiO₂ ultrathin layers not only facilitated CO₂ capture from air, but also benefited for the CO₂ activation process and the following rate-limiting COOH* scission step, which definitely promoted CO₂ photoreduction performance in the fresh air.

In summary, for directly capturing and effectively activating atmospheric CO₂ molecules, abundant Ti³⁺-V_o-Ti³⁺ defective sites confined in TiO₂ ultrathin layers were designed and fabricated. XPS spectra demonstrated the presence of Ti³⁺ and V_o in the synthetic TiO₂ ultrathin layers, while Bader charge calculations and PAL spectra first disclosed the formed oxygen vacancies were associated with the formation of two adjacent Ti³⁺ sites. *In-situ* FTIR spectra and DFT calculations unveiled these generated Ti³⁺-V_o-Ti³⁺ sites could simultaneously bond with the C and O atoms of CO₂ molecules to form a stable bridge adsorption configuration, which converted the endoergic chemisorption step to an exoergic process, showing feasibility for capturing CO₂ directly from the air. More importantly, the two adjacent Ti³⁺ sites not only facilitated CO₂ activation into COOH* through forming a stable Ti³⁺-C-O-Ti³⁺ configuration, but also lowered the rate-limiting energy barrier from 0.75 to 0.45 eV *via* stabilizing the produced CO*. As a result, the Ti³⁺-V_o-TiO₂ ultrathin layers could directly capture CO₂ in the fresh air and selectively photofix the atmospheric CO₂ into near-unity CO, in which the CO₂-to-CO conversion ratio was ca. 20.2%, giving possibilities of the direct implementation of CO₂ in the fresh air. Thus, this defect design concept for tailoring atmospheric CO₂ capture and activation opens up broad possibilities for “real-world” application of solar fuel production.

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Conflict of interest The authors declare no conflict of interest.

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