

• Last time:

• Gauss's law:

$$\frac{\partial R^{\delta}_{\alpha\beta}}{\partial u^\gamma} - \frac{\partial R^{\delta}_{\alpha\gamma}}{\partial u^\beta} + R^{\rho}_{\alpha\beta} R^{\delta}_{\rho\gamma} - R^{\rho}_{\alpha\gamma} R^{\delta}_{\rho\beta}$$

$$= b_{\alpha\beta} b^{\delta}_\gamma - b_{\alpha\gamma} b^{\delta}_\beta$$

• Codazzi's law:

$$\frac{\partial b_{\alpha\beta}}{\partial u^\gamma} - \frac{\partial b_{\alpha\gamma}}{\partial u^\beta} = R^{\delta}_{\alpha\gamma} b_{\delta\beta} - R^{\delta}_{\alpha\beta} b_{\delta\gamma}$$

??: $R^{\delta}_{\alpha\beta\gamma} :=$

$$\frac{\partial R^{\delta}_{\alpha\beta}}{\partial u^\gamma} - \frac{\partial R^{\delta}_{\alpha\gamma}}{\partial u^\beta} + R^{\rho}_{\alpha\beta} R^{\delta}_{\rho\gamma} - R^{\rho}_{\alpha\gamma} R^{\delta}_{\rho\beta}$$

Q1. $R^{\delta}_{\alpha\beta\gamma}$ 是什么量?

$R^{\alpha}_{\beta\gamma}$ 是什么量?

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Q2. 对于?

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To answer these two questions we introduce tensor analysis.

• Brief introduction:

We call

$$\overline{T}^{\alpha_1 \dots \alpha_r}_{\beta_1 \dots \beta_s}(u^1, u^2)$$

a tensor of type (r, s)

if under the change of
coords $(u^1, u^2) \mapsto (\tilde{u}^1, \tilde{u}^2)$

$\overline{T}^{\alpha_1 \dots \alpha_r}_{\beta_1 \dots \beta_s}$ transform as follows:

$$\tilde{T}_{\alpha_1 \dots \alpha_r}^{\beta_1 \dots \beta_s}(u^1, u^2)$$

$$= T_{\alpha_1 \dots \alpha_r}^{\beta_1 \dots \beta_s} \frac{\partial u^{\beta_1}}{\partial u^{\alpha_1}} \dots \frac{\partial u^{\beta_s}}{\partial u^{\alpha_r}} \cdot \frac{\tilde{u}^{\alpha_1}}{\partial u^{\alpha_1}} \dots \frac{\tilde{u}^{\alpha_r}}{\partial u^{\alpha_r}}$$

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Therefore:

$$\tilde{T}_{\alpha_1 \dots \alpha_r}^{\beta_1 \dots \beta_s}(u^1, u^2) du^{\beta_1} \otimes \dots \otimes du^{\beta_s} \frac{\partial}{\partial u^{\alpha_1}} \otimes \dots \otimes \frac{\partial}{\partial u^{\alpha_r}}$$

is invariant under coords transf.

上 $\alpha_1, \dots, \alpha_r$ 称为 反变的

下 β_1, \dots, β_s --- 扭变的.

Lemma. 1

$R^{\delta}_{\alpha\beta\gamma}$ 为具有一个反变分量和三个扭变分量
的 $(1,3)$ -型 张量.

为证这个引理, 我们先看 $g_{\alpha\beta}, g^{\alpha\beta}, \tilde{g}^{\alpha\beta}$ 在
坐标变换下如何改变.

Lemma. 2 $\tilde{g}^{\alpha\beta}$ 不是张量.

pf. Recall that

$$\tilde{g}^{\alpha\beta} = \frac{1}{2} \left(\frac{\partial \tilde{g}^{\alpha\beta}}{\partial u^\rho} + \frac{\partial \tilde{g}^{\beta\alpha}}{\partial u^\rho} - \frac{\partial \tilde{g}^{\alpha\beta}}{\partial u^\lambda} \right) \tilde{g}^{\lambda\rho}$$

$$\tilde{g}^{\alpha\beta} = \frac{1}{2} \left(\frac{\partial \tilde{g}^{\alpha\beta}}{\partial \tilde{u}^\rho} + \frac{\partial \tilde{g}^{\beta\alpha}}{\partial \tilde{u}^\rho} - \frac{\partial \tilde{g}^{\alpha\beta}}{\partial \tilde{u}^\lambda} \right) \tilde{g}^{\lambda\rho}$$

$$= \frac{1}{2} \left(\frac{\partial}{\partial \tilde{u}^\rho} \left(g_{\rho\sigma} \frac{\partial u^\rho}{\partial \tilde{u}^\alpha} \frac{\partial u^\sigma}{\partial \tilde{u}^\beta} \right) \right)$$

$$+ \frac{\partial}{\partial \tilde{u}^\lambda} \left(g_{\rho\sigma} \frac{\partial u^\rho}{\partial \tilde{u}^\beta} \frac{\partial u^\sigma}{\partial \tilde{u}^\alpha} \right)$$

$$- \frac{\partial}{\partial \tilde{u}^\beta} \left(g_{\rho\sigma} \frac{\partial u^\rho}{\partial \tilde{u}^\alpha} \frac{\partial u^\sigma}{\partial \tilde{u}^\lambda} \right) \right) \tilde{g}^{\lambda\rho}$$

$$\begin{aligned}
&= \frac{1}{2} \left(\frac{\partial g_{\rho\sigma}}{\partial u^\beta} \frac{\partial u^\rho}{\partial u^\alpha} \frac{\partial u^\sigma}{\partial u^\beta} + g_{\rho\sigma} \frac{\partial u^\rho}{\partial u^\alpha} \frac{\partial u^\sigma}{\partial u^\beta} \right) \textcircled{28} \\
&\quad + g_{\rho\sigma} \frac{\partial u^\rho}{\partial u^\alpha} \frac{\partial^2 u^\sigma}{\partial u^\beta \partial u^\alpha} \\
&+ \frac{\partial g_{\rho\sigma}}{\partial u^\alpha} \frac{\partial u^\rho}{\partial u^\beta} \frac{\partial u^\sigma}{\partial u^\beta} + g_{\rho\sigma} \frac{\partial^2 u^\rho}{\partial u^\beta \partial u^\alpha} \frac{\partial u^\sigma}{\partial u^\beta} \\
&\quad + g_{\rho\sigma} \frac{\partial u^\rho}{\partial u^\beta} \frac{\partial^2 u^\sigma}{\partial u^\beta \partial u^\alpha} \\
&- \frac{\partial g_{\rho\sigma}}{\partial u^\beta} \frac{\partial u^\rho}{\partial u^\alpha} \frac{\partial u^\sigma}{\partial u^\beta} - g_{\rho\sigma} \frac{\partial^2 u^\rho}{\partial u^\beta \partial u^\alpha} \frac{\partial u^\sigma}{\partial u^\beta} \\
&\quad - g_{\rho\sigma} \frac{\partial u^\rho}{\partial u^\alpha} \frac{\partial^2 u^\sigma}{\partial u^\beta \partial u^\alpha} \right) \tilde{g}^{33}
\end{aligned}$$

$$\begin{aligned}
&= \frac{1}{2} \left(\frac{\partial g_{\rho\sigma}}{\partial u^\beta} \frac{\partial u^\rho}{\partial u^\alpha} \frac{\partial u^\sigma}{\partial u^\beta} \underbrace{g_{\phi\psi}}_{\text{---}} \frac{\partial u^\phi}{\partial u^\alpha} \frac{\partial u^\psi}{\partial u^\beta} \right. \\
&+ \frac{\partial g_{\rho\sigma}}{\partial u^\beta} \frac{\partial u^\rho}{\partial u^\alpha} \frac{\partial u^\sigma}{\partial u^\beta} \underbrace{g_{\phi\psi}}_{\text{---}} \frac{\partial u^\phi}{\partial u^\alpha} \frac{\partial u^\psi}{\partial u^\beta} \\
&- \frac{\partial g_{\rho\sigma}}{\partial u^\beta} \frac{\partial u^\rho}{\partial u^\alpha} \frac{\partial u^\sigma}{\partial u^\beta} \underbrace{g_{\phi\psi}}_{\text{---}} \frac{\partial u^\phi}{\partial u^\alpha} \frac{\partial u^\psi}{\partial u^\beta} \left. \right)
\end{aligned}$$

$$\begin{aligned}
&+ 2 g_{\rho\sigma} \frac{\partial u^\rho}{\partial u^\alpha} \frac{\partial u^\sigma}{\partial u^\beta} \tilde{g}^{33} \textcircled{28} \\
&= \frac{1}{2} \left(\frac{\partial g_{\rho\phi}}{\partial u^\beta} \frac{\partial u^\rho}{\partial u^\alpha} \frac{\partial u^\phi}{\partial u^\beta} g_{\phi\psi} + \frac{\partial u^\delta}{\partial u^\beta} \right. \\
&+ \frac{\partial g_{\rho\phi}}{\partial u^\beta} \frac{\partial u^\rho}{\partial u^\alpha} \frac{\partial u^\phi}{\partial u^\beta} g_{\phi\psi} \frac{\partial u^\delta}{\partial u^\beta} \\
&- \frac{\partial g_{\rho\phi}}{\partial u^\beta} \frac{\partial u^\rho}{\partial u^\alpha} \frac{\partial u^\phi}{\partial u^\beta} g_{\phi\psi} \frac{\partial u^\delta}{\partial u^\beta} \left. \right) \\
&+ g_{\rho\sigma} \frac{\partial^2 u^\rho}{\partial u^\beta \partial u^\alpha} \frac{\partial u^\sigma}{\partial u^\beta} g^{\phi\psi} \frac{\partial u^\delta}{\partial u^\alpha} \\
&= \frac{1}{2} \left(\frac{\partial g_{\rho\phi}}{\partial u^\beta} \frac{\partial u^\rho}{\partial u^\alpha} \frac{\partial u^\phi}{\partial u^\beta} g_{\phi\psi} \frac{\partial u^\delta}{\partial u^\beta} \right. \\
&+ \frac{\partial g_{\phi\psi}}{\partial u^\beta} \frac{\partial u^\rho}{\partial u^\alpha} \frac{\partial u^\phi}{\partial u^\beta} g_{\phi\psi} \frac{\partial u^\delta}{\partial u^\alpha} \\
&- \frac{\partial g_{\rho\phi}}{\partial u^\beta} \frac{\partial u^\rho}{\partial u^\alpha} \frac{\partial u^\phi}{\partial u^\beta} g_{\phi\psi} \frac{\partial u^\delta}{\partial u^\alpha} \left. \right)
\end{aligned}$$

$$+ g_{\rho\sigma} \frac{\partial^2 u^\rho}{\partial u^\alpha \partial u^\beta} g^{\sigma\gamma} \frac{\partial u^\delta}{\partial u^\gamma}$$

$$= P_{\rho\varepsilon}^{\gamma} \frac{\partial u^\varepsilon}{\partial u^\beta} \frac{\partial u^\rho}{\partial u^\alpha} \frac{\partial u^\delta}{\partial u^\gamma} + \boxed{\frac{\partial^2 u^\rho}{\partial u^\alpha \partial u^\beta} \frac{\partial u^\delta}{\partial u^\gamma}}$$

this term in
general is not zero

Therefore,

$P_{\rho\beta}^\gamma$ is not a
tensor. #

If of Lemma 1.

$$R_{\alpha\beta\gamma}^\delta :=$$

$$\frac{\partial P_{\alpha\beta}^\delta}{\partial u^\gamma} - \frac{\partial P_{\alpha\gamma}^\delta}{\partial u^\beta} + P_{\alpha\beta}^\rho P_{\rho\gamma}^\delta - P_{\alpha\gamma}^\rho P_{\rho\beta}^\delta$$

if.

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$$\begin{aligned} \tilde{R}_{\alpha\beta\gamma}^\delta &= \frac{\partial \tilde{P}_{\alpha\beta}^\delta}{\partial u^\gamma} - \frac{\partial \tilde{P}_{\alpha\gamma}^\delta}{\partial u^\beta} + \tilde{P}_{\alpha\beta}^\rho \tilde{P}_{\rho\gamma}^\delta - \tilde{P}_{\alpha\gamma}^\rho \tilde{P}_{\rho\beta}^\delta \\ &= \frac{\partial}{\partial u^\gamma} \left(P_{\rho\varepsilon}^{\gamma} \frac{\partial u^\varepsilon}{\partial u^\beta} \frac{\partial u^\rho}{\partial u^\alpha} \frac{\partial u^\delta}{\partial u^\gamma} + \frac{\partial^2 u^\rho}{\partial u^\alpha \partial u^\beta} \frac{\partial u^\delta}{\partial u^\gamma} \right) \frac{\partial u^\phi}{\partial u^\gamma} \\ &\quad - \frac{\partial}{\partial u^\beta} \left(P_{\rho\varepsilon}^{\gamma} \frac{\partial u^\varepsilon}{\partial u^\gamma} \frac{\partial u^\rho}{\partial u^\alpha} \frac{\partial u^\delta}{\partial u^\beta} + \frac{\partial^2 u^\rho}{\partial u^\alpha \partial u^\gamma} \frac{\partial u^\delta}{\partial u^\beta} \right) \frac{\partial u^\phi}{\partial u^\beta} \\ &\quad + \left(P_{\mu\varepsilon}^{\gamma} \frac{\partial u^\varepsilon}{\partial u^\beta} \frac{\partial u^\mu}{\partial u^\alpha} \frac{\partial u^\rho}{\partial u^\gamma} + \frac{\partial^2 u^\mu}{\partial u^\alpha \partial u^\beta} \frac{\partial u^\rho}{\partial u^\gamma} \right) \\ &\quad \left(P_{\mu\varepsilon}^{\gamma} \frac{\partial u^\varepsilon}{\partial u^\gamma} \frac{\partial u^\mu}{\partial u^\beta} \frac{\partial u^\delta}{\partial u^\alpha} + \frac{\partial^2 u^\mu}{\partial u^\beta \partial u^\gamma} \frac{\partial u^\delta}{\partial u^\alpha} \right) \\ &\quad - \left(P_{\mu\varepsilon}^{\gamma} \frac{\partial u^\varepsilon}{\partial u^\gamma} \frac{\partial u^\mu}{\partial u^\alpha} \frac{\partial u^\rho}{\partial u^\delta} + \frac{\partial^2 u^\mu}{\partial u^\alpha \partial u^\gamma} \frac{\partial u^\rho}{\partial u^\delta} \right) \\ &\quad \left(P_{\mu\varepsilon}^{\gamma} \frac{\partial u^\varepsilon}{\partial u^\beta} \frac{\partial u^\mu}{\partial u^\rho} \frac{\partial u^\delta}{\partial u^\alpha} + \frac{\partial^2 u^\mu}{\partial u^\beta \partial u^\rho} \frac{\partial u^\delta}{\partial u^\alpha} \right) \end{aligned} \quad (80)$$

$$\begin{aligned}
 &= \frac{\partial \Gamma_{\mu\nu}^A}{\partial u^\phi} \frac{\partial h^3}{\partial u^\phi} \frac{\partial h^1}{\partial u^\phi} \frac{\partial h^2}{\partial u^\phi} \frac{\partial u^\phi}{\partial u^\phi} \\
 &+ \overbrace{\Gamma_{\mu\nu}^A \frac{\partial h^3}{\partial u^\phi} \frac{\partial h^1}{\partial u^\phi} \frac{\partial h^2}{\partial u^\phi} + \Gamma_{\mu\nu}^A \frac{\partial h^3}{\partial u^\phi} \frac{\partial h^2}{\partial u^\phi} \frac{\partial h^1}{\partial u^\phi}}^{\text{blue line}} \\
 &+ \overbrace{\Gamma_{\mu\nu}^A \frac{\partial h^3}{\partial u^\phi} \frac{\partial h^1}{\partial u^\phi} \frac{\partial h^2}{\partial u^\phi} + \Gamma_{\mu\nu}^A \frac{\partial h^3}{\partial u^\phi} \frac{\partial h^2}{\partial u^\phi} \frac{\partial h^1}{\partial u^\phi}}^{\text{green line}} \\
 &+ \overbrace{\Gamma_{\mu\nu}^A \frac{\partial h^3}{\partial u^\phi} \frac{\partial h^1}{\partial u^\phi} \frac{\partial h^2}{\partial u^\phi} + \Gamma_{\mu\nu}^A \frac{\partial h^3}{\partial u^\phi} \frac{\partial h^2}{\partial u^\phi} \frac{\partial h^1}{\partial u^\phi}}^{\text{red line}} \\
 &+ \overbrace{\left(\Gamma_{\mu\nu}^A \frac{\partial h^3}{\partial u^\phi} \frac{\partial h^1}{\partial u^\phi} \frac{\partial h^2}{\partial u^\phi} + \Gamma_{\mu\nu}^A \frac{\partial h^3}{\partial u^\phi} \frac{\partial h^2}{\partial u^\phi} \frac{\partial h^1}{\partial u^\phi} \right) \left(\Gamma_{\mu\nu}^A \frac{\partial h^3}{\partial u^\phi} \frac{\partial h^1}{\partial u^\phi} \frac{\partial h^2}{\partial u^\phi} + \Gamma_{\mu\nu}^A \frac{\partial h^3}{\partial u^\phi} \frac{\partial h^2}{\partial u^\phi} \frac{\partial h^1}{\partial u^\phi} \right)}^{\text{brown line}}
 \end{aligned}$$

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$$\begin{aligned}
 & - \left(P^4 \frac{\partial u^M}{\partial n^\alpha} \frac{\partial u^\beta}{\partial n^\gamma} \frac{\partial u^\delta}{\partial n^\epsilon} \right) \left(P^4 \frac{\partial u^\epsilon}{\partial n^\beta} \frac{\partial u^M}{\partial n^\delta} \frac{\partial u^\gamma}{\partial n^\alpha} \right) \\
 & + \underbrace{\frac{\partial^2 u^i}{\partial n^\alpha \partial n^\beta} \frac{\partial u^\rho}{\partial n^\epsilon} \left(P^4 \frac{\partial u^\delta}{\partial n^\beta} \frac{\partial u^M}{\partial n^\rho} \frac{\partial u^\gamma}{\partial n^\alpha} + \frac{\partial^2 u^i}{\partial n^\beta \partial n^\gamma} \frac{\partial u^\delta}{\partial n^\epsilon} \right)}_{\text{blue line}} \\
 & + \underbrace{\frac{\partial^2 u^i}{\partial n^\beta \partial n^\gamma} \frac{\partial u^\delta}{\partial n^\epsilon} \left(P^4 \frac{\partial u^\delta}{\partial n^\beta} \frac{\partial u^M}{\partial n^\rho} \frac{\partial u^\gamma}{\partial n^\alpha} \right)}_{\text{green line}} \\
 & - \underbrace{\frac{\partial^2 u^i}{\partial n^\alpha \partial n^\beta} \frac{\partial u^\rho}{\partial n^\epsilon} \left(P^4 \frac{\partial u^\delta}{\partial n^\beta} \frac{\partial u^M}{\partial n^\rho} \frac{\partial u^\gamma}{\partial n^\alpha} + \frac{\partial^2 u^i}{\partial n^\beta \partial n^\gamma} \frac{\partial u^\delta}{\partial n^\epsilon} \right)}_{\text{orange line}} \\
 & - \underbrace{\frac{\partial^2 u^i}{\partial n^\beta \partial n^\gamma} \frac{\partial u^\delta}{\partial n^\epsilon} \left(P^4 \frac{\partial u^\delta}{\partial n^\beta} \frac{\partial u^M}{\partial n^\rho} \frac{\partial u^\gamma}{\partial n^\alpha} \right)}_{\text{red line}} \\
 & = \text{---}
 \end{aligned}$$

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$$\begin{aligned}
& + \frac{\partial u^4}{\partial u^\alpha \partial u^\beta} \left(\frac{\partial^2 u^\delta}{\partial u^\alpha \partial u^\beta} \frac{\partial u^\phi}{\partial u^\delta} + \frac{\partial u^\rho}{\partial u^\alpha} \frac{\partial^2 u^\delta}{\partial u^\rho \partial u^\beta} \frac{\partial u^\phi}{\partial u^\delta} \right) \quad (85) \\
& - \frac{\partial^2 u^\delta}{\partial u^\alpha \partial u^\beta} \left(\frac{\partial^2 u^\delta}{\partial u^\alpha \partial u^\beta} \frac{\partial u^\phi}{\partial u^\delta} + \frac{\partial u^\rho}{\partial u^\alpha} \frac{\partial^2 u^\delta}{\partial u^\rho \partial u^\beta} \frac{\partial u^\phi}{\partial u^\delta} \right) \\
= & 0 - \Gamma^\gamma_{\mu\delta} \frac{\partial u^\mu}{\partial u^\alpha} \frac{\partial u^\delta}{\partial u^\gamma} \frac{\partial}{\partial u^\beta} \left(\frac{\partial u^\delta}{\partial u^\alpha} \frac{\partial u^\phi}{\partial u^\beta} \right) \\
& + \Gamma^\gamma_{\mu\delta} \frac{\partial u^\mu}{\partial u^\alpha} \frac{\partial u^\delta}{\partial u^\beta} \frac{\partial}{\partial u^\gamma} \left(\frac{\partial u^\delta}{\partial u^\alpha} \frac{\partial u^\phi}{\partial u^\beta} \right) \\
& + \frac{\partial^2 u^\delta}{\partial u^\alpha \partial u^\beta} \frac{\partial}{\partial u^\delta} \left(\frac{\partial u^\delta}{\partial u^\alpha} \frac{\partial u^\phi}{\partial u^\beta} \right) \\
& - \frac{\partial^2 u^\delta}{\partial u^\alpha \partial u^\beta} \frac{\partial}{\partial u^\delta} \left(\frac{\partial u^\delta}{\partial u^\beta} \frac{\partial u^\phi}{\partial u^\alpha} \right) \\
= & 0. \quad \#.
\end{aligned}$$

• 引理

例. 由 $\Gamma^\gamma_{\mu\delta}$ 的定义知

对称性

$$R_{\alpha\beta\gamma\delta} = R_{\alpha\beta\delta\gamma} = -R_{\alpha\delta\beta\gamma} = -R_{\alpha\gamma\beta\delta}$$

§3.2. 刚体运动定理

Thm. 若 S_1, S_2 为定义在 D 上的两个正则曲面
 $\vec{r}(u, v), \vec{r}(u, v)$. 若 $\forall (u, v) \in D$
 S_1, S_2 在 (u, v) 处有相同的 I 和 II, 则
 在 \mathbb{R}^3 中的刚体运动 T 及

$$\vec{r} = T \circ \vec{r}.$$

pf. 任取 $\exists u_0 = (u_0^1, u_0^2) \in D$

$$\left. \begin{array}{l} I_1|_{u_0} = I_2|_{u_0} \\ II_1|_{u_0} = II_2|_{u_0} \end{array} \right\} \Rightarrow \text{找 } \left\{ \vec{r}(u_0), \vec{r}_u(u_0), \vec{n}(u_0) \right\}$$

经过一个适当的刚体运动后, 与找

$$\left\{ \tilde{\vec{r}}(u_0), \tilde{\vec{r}}_u(u_0), \tilde{\vec{n}}(u_0) \right\} \text{ 相合.}$$

$$\left. \begin{array}{l} \tilde{\vec{r}}(u) = T \circ \vec{r}(u_0) \\ \tilde{\vec{r}}_u(u) = T \circ \vec{r}_{u_0}(u_0) \\ \tilde{\vec{n}}(u) = T \circ \vec{n}(u_0) \end{array} \right\} \quad (1)$$

自然转置
由于运动方程的系数由 I 与 II 完全确定.

而 $\frac{I_1}{II_1} = \frac{I_2}{II_2}$, 由一阶线性 PDE 方程组初值问题

解的唯一性及①知: $\tilde{r} \equiv T \cdot r$. #

存在性定理

Thm. 对于 $I = g_{\alpha\beta} du^\alpha du^\beta$

$$II = b_{\alpha\beta} du^\alpha du^\beta, \quad u = (u^1, u^2) \in D$$

其中 D 为 \mathbb{R}^2 上的单连通区域, $g_{11}, g_{12} = g_{21}, g_{22}$,
 $b_{11}, b_{12} = b_{21}, b_{22}$ 为 D 上给定的光滑函数, 且 $g_{11} g_{22} - g_{12}^2 > 0$.

定义: $\tilde{P}_{\alpha\beta}^\gamma = \frac{1}{2} g^{\gamma\delta} \left(\frac{\partial g_{\alpha\delta}}{\partial u^\beta} + \frac{\partial g_{\beta\delta}}{\partial u^\alpha} - \frac{\partial g_{\alpha\beta}}{\partial u^\delta} \right)$ (Gauss 方程)

若 $\tilde{P}_{\alpha\beta}^\gamma, b_{\alpha\beta}$ 满足高斯方程和 Codazzi 方程,

则 $\forall u \in D$, 存在 u 的一个邻域 $U \subseteq D$, 以及
定义在 U 上的曲面 $\tilde{r}: U \rightarrow \mathbb{R}^3$, 使.

I, II 为其第一、第二基本形式.

Wf. Consider the following linear PDE system:

$$\begin{cases} \frac{\partial \vec{p}}{\partial u^\alpha} = \vec{P}_\alpha, \\ \frac{\partial \vec{P}_\alpha}{\partial u^\beta} = \tilde{P}_{\alpha\beta}^\gamma \vec{P}_\gamma + b_{\alpha\beta} \vec{f}, \\ \frac{\partial \vec{f}}{\partial u^\alpha} = -b_\alpha^\gamma \vec{P}_\gamma \end{cases} \quad (*)$$

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这个一阶变系数 线性方程组 的初值问题条件为:

$$\partial_{u^\beta} (\partial_{u^\alpha} \vec{p}) = \alpha \leftrightarrow \beta$$

$$\partial_{u^\beta} (\partial_{u^\alpha} \vec{P}_\alpha) = \beta \leftrightarrow \gamma$$

$$\partial_{u^\beta} (\partial_{u^\alpha} \vec{f}) = \alpha \leftrightarrow \beta$$

$$\Leftrightarrow \begin{cases} \tilde{P}_{\alpha\beta}^\gamma = \tilde{P}_{\beta\alpha}^\gamma, \\ b_{\alpha\beta} = b_{\beta\alpha}, \\ \text{Gauss 方程以及 Codazzi 方程.} \end{cases} \quad (*)$$

对 $u_0 \in D$, 任取向量 $\vec{P}_{\alpha_0}, \vec{f}_0$ 满足:

$$\langle \vec{f}_0, \vec{f}_0 \rangle = 1,$$

$$\langle \vec{P}_{\alpha_0}, \vec{f}_0 \rangle = 0,$$

$$\langle \vec{P}_{\alpha_0}, \vec{P}_{\beta_0} \rangle = \tilde{P}_{\alpha_0\beta_0}(u_0).$$

$$(\vec{P}_{\alpha_1}, \vec{P}_{\alpha_2}, \vec{f}_0) > 0$$

Consider (*) 的初值问题:

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$$\vec{p}(u_0) = \vec{p}_0 \quad (\text{假设 } \mathbb{R}^3 \text{ 中一式})$$

$$\vec{p}_\alpha(u_0) = \vec{p}_{0\alpha}$$

$$\vec{f}(u_0) = \vec{f}_0$$

由于(*)成立 \Rightarrow 存在 u_0 的一个邻域 $U \subseteq D$, s.t.
(*) 在 U 中有解:

$\vec{p}, \vec{p}_\alpha, \vec{f} \Rightarrow \vec{p}$ 给出 $U \rightarrow \mathbb{R}^3$ 的一个
参数曲面.

$$\begin{aligned} (*) \Rightarrow \frac{\partial \langle \vec{p}_\alpha, \vec{p}_\beta \rangle}{\partial u^\gamma} &= P_{\alpha\gamma}^3 \langle \vec{p}_3, \vec{p}_\beta \rangle \\ &\quad + P_{\beta\gamma}^3 \langle \vec{p}_3, \vec{p}_\alpha \rangle \\ &\quad + b_{\alpha\gamma} \langle \vec{p}_\beta, \vec{f} \rangle \\ &\quad + b_{\beta\gamma} \langle \vec{p}_\alpha, \vec{f} \rangle \\ \frac{\partial \langle \vec{p}_\alpha, \vec{f} \rangle}{\partial u^\beta} &= P_{\alpha\beta}^3 \langle \vec{p}_3, \vec{f} \rangle \\ &\quad + b_{\alpha\beta} \langle \vec{f}, \vec{f} \rangle \\ &\quad - b_\beta^3 \langle \vec{p}_\alpha, \vec{p}_3 \rangle \end{aligned}$$

$$\frac{\partial \langle \vec{f}, \vec{f} \rangle}{\partial u^\alpha} = -2b_\alpha^\beta \langle \vec{p}_\beta, \vec{f} \rangle \quad (88)$$

初值条件

$$\langle \vec{p}_\alpha, \vec{p}_\beta \rangle, \langle \vec{p}_\alpha, \vec{f} \rangle, \langle \vec{f}, \vec{f} \rangle$$

而 -阶 线 性 方 程 组 .

显然 $g_{\alpha\beta}, 0, 1$ 为一组解.

给定初值 $g_{\alpha\beta}(u_0), 0, 1$. 解之 -.

$$\begin{aligned} \text{由于 } \langle \vec{p}_\alpha, \vec{p}_\beta \rangle(u_0) &= g_{\alpha\beta}(u_0), \\ \langle \vec{p}_\alpha, \vec{f} \rangle(u_0) &= 0, \\ \langle \vec{f}, \vec{f} \rangle(u_0) &= 1 \end{aligned}$$

$$\begin{aligned} \Rightarrow \langle \vec{p}_\alpha, \vec{p}_\beta \rangle &\equiv g_{\alpha\beta} \\ \langle \vec{p}_\alpha, \vec{f} \rangle &\equiv 0 \\ \langle \vec{f}, \vec{f} \rangle &\equiv 1. \end{aligned}$$

$\Rightarrow \vec{p}_1, \vec{p}_2$ 在 U 上处处线性无关.

且由 $\vec{p}_3 = \partial_{u^\alpha} \vec{p}$ 知: \vec{p} 以 I 为其第一基本形式.
 $\vec{p}_1, \vec{p}_2, \vec{f}$ 为右手, $\langle \vec{f}, \vec{f} \rangle \equiv 1$ 及

$$\frac{\partial \vec{P}_\alpha}{\partial u^\beta} = \vec{P}_{\alpha\beta}^{\gamma} \vec{P}_\gamma + b_{\alpha\beta} \vec{f}$$

(3)

知: \vec{P} 为曲面 S 的第二基本形式. 定义 $\vec{P} = \vec{P}_{\alpha\beta}^{\gamma} du^\alpha du^\beta$.

作业: ① 给出 $R_{\alpha\beta\gamma\delta}$ 对称性的内蕴证明.

② $P_{102} : 1$

$P_{103} : 3, 5, 9$

§4. 联络与测地线

§4.1 平行移动

S : 光滑曲面. 正则. $\vec{r} = \vec{r}(u^1, u^2)$ $\vec{r}: D \rightarrow \mathbb{R}^3$
设 $\vec{X}(u^1, u^2)$ 为 S 上的一个切向量场. $(u^1, u^2) \mapsto \vec{X}(u^1, u^2)$.
 $\vec{X} = X^\alpha(u^1, u^2) \vec{r}_{u^\alpha}(u^1, u^2)$.

若 \vec{X} 是光滑的, if x^1, x^2 均是光滑的. (on D)

$$d\vec{X} = dX^\alpha \vec{r}_{u^\alpha} + X^\alpha d\vec{r}_{u^\alpha}$$

$$\begin{aligned} &= dX^\alpha \vec{r}_{u^\alpha} + X^\alpha \left(\vec{P}_{\alpha\beta}^{\gamma} du^\beta \vec{r}_{u^\alpha} + b_{\alpha\beta} du^\beta \vec{n} \right) \\ &= (dX^\alpha + X^\beta \vec{P}_{\beta\alpha}^{\alpha} du^\beta) \vec{r}_{u^\alpha} + X^\alpha b_{\alpha\beta} du^\beta \vec{n}. \end{aligned}$$

定义: $D\vec{X}(u^1, u^2) := (dX^\alpha + X^\beta \vec{P}_{\beta\alpha}^{\alpha} du^\beta) \vec{r}_{u^\alpha}$.

称 $D\vec{X}$ 为曲面 S 上的切向量场 \vec{X} 的平行微分.

(M上, $\forall p \in S$, $D\vec{X}(p)$ 为 $d\vec{X}(p)$ 向 $T_p S$ 的正交投影).

$$\text{即: } D\vec{X} = dX^\alpha + X^\beta \vec{P}_{\beta\alpha}^{\alpha} du^\beta.$$

记:

$$D\vec{X} = D\vec{X}^\alpha \vec{r}_{u^\alpha}$$

称 $D\vec{X}$ 为 \vec{X} 在 S 上的切变微分. $D\vec{X} = \frac{\partial \vec{X}}{\partial u^\alpha} du^\alpha + X^\beta \vec{P}_{\beta\alpha}^{\alpha} du^\beta$
Thm. $\forall \vec{x}, \vec{y} \in \mathcal{C}(S)$, 有: $= \left(\frac{\partial \vec{X}}{\partial u^\alpha} + X^\beta \vec{P}_{\beta\alpha}^{\alpha} \right) du^\alpha$

$$\textcircled{1} D(\vec{x} + \vec{y}) = D(\vec{x}) + D(\vec{y})$$

$$\textcircled{2} D(f \cdot \vec{x}) = df \cdot \vec{x} + f \cdot D\vec{x}, \quad \forall f \in C^\infty(D)$$

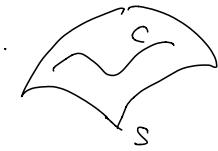
$$\textcircled{3} d\langle \vec{x}, \vec{y} \rangle = \langle D\vec{x}, \vec{y} \rangle + \langle \vec{x}, D\vec{y} \rangle$$

习题 exercise. 未

设 $C: u^\alpha = u^\alpha(t)$ 为 S 上的一条曲线.

设 $\vec{X}(t)$ 为 S 上沿 C 的一个切向量场

$\frac{d\vec{X}(t)}{dt}$ 为 \mathbb{R}^3 中沿 C 的一个向量场.



一般来说, $\frac{d\vec{X}(t)}{dt}$ 不是 S 上的切向量场.

定义: $\frac{D\vec{X}(t)}{dt} := \frac{d\vec{X}(t)}{dt}$ 在 $T_{(u^{1(t)}, u^{2(t)})} S$ 上的投影.

$$\vec{X}(t) = X^{(t)} \vec{r}_{u^\alpha}(u^{1(t)}, u^{2(t)})$$

$$\frac{d\vec{X}(t)}{dt} = \frac{dX^{(t)}}{dt} \vec{r}_{u^\alpha} + X^{(t)} \frac{\partial \vec{r}_{u^\alpha}}{\partial u^\alpha} \frac{du^{(t)}}{dt}$$

$$= \underbrace{\left(\frac{dX^{(t)}}{dt} + \vec{P}_{\beta\alpha}^{\alpha} X^{(t)} \frac{du^{(t)}}{dt} \right)}_{\text{括号部分}} \vec{r}_{u^\alpha} + b_{\alpha\beta} X^{(t)} \frac{du^{(t)}}{dt} \vec{n}$$

$$\Rightarrow \frac{D\vec{x}(t)}{dt} = \left(\frac{dx^\alpha}{dt} + \Gamma_{\beta\gamma}^\alpha x^\beta \frac{du^\gamma(t)}{dt} \right) \vec{r}_\omega \quad (9)$$

若 $\frac{D\vec{x}(t)}{dt}$ 为 S 上沿 C 定义的 S 上的切向量场 X

沿 C 在 S 上的切变导数.

$$\frac{dx^\alpha}{dt} + \Gamma_{\beta\gamma}^\alpha x^\beta \frac{du^\gamma(t)}{dt} =: \frac{Dx^\alpha}{dt} \text{ 为分量 } X^\alpha \text{ 沿 } C \text{ 的切变导数.}$$

Recall that $Dx^\alpha = \left(\frac{\partial x^\alpha}{\partial u^\gamma} + x^\beta \Gamma_{\beta\gamma}^\alpha \right) du^\gamma$

$\frac{\partial x^\alpha}{\partial u^\gamma} + x^\beta \Gamma_{\beta\gamma}^\alpha$ 为沿参数曲线的切变导数.

记为 X^α_γ . 因而, $Dx^\alpha = X^\alpha_\gamma du^\gamma$.

定义: 设 $\vec{Y}(t)$ 是 S 上沿 $C: u^\gamma = u^\gamma(t)$ 定义的 S 上的光滑切向量场.

if $\frac{D\vec{Y}(t)}{dt} = 0$, 则称 $\vec{Y}(t)$ 在 C 在 S 上平行的.

$$\cdot \frac{D\vec{X}(t)}{dt} = 0 \Leftrightarrow \frac{dX^\alpha}{dt} + \Gamma_{\beta\gamma}^\alpha X^\beta \frac{du^\gamma}{dt} = 0. \quad (*)$$

$$C: u^\gamma = u^\gamma(t), \quad t \in [c, d] \subseteq [0, b]$$

$$X|_{t=t_0} = X_0$$

ODE 理论 $\Rightarrow \vec{X}(t)$ 为 $(*)$ 的唯一解.

若 $\vec{X}(t) := X^\alpha(t) \vec{r}_\omega$ 为 $X^\alpha(\vec{r}_\omega(u^\gamma(t)), u^\gamma(t))$ 沿 C 的平行场.

若 $\vec{X}(t)$ 沿 C 在 S 上平行,

$$\frac{d}{dt} \langle \vec{X}(t), \vec{Y}(t) \rangle = \langle \frac{D\vec{X}(t)}{dt}, \vec{Y}(t) \rangle + \langle \vec{X}(t), \frac{D\vec{Y}(t)}{dt} \rangle$$

$\stackrel{(8.6)}{\Rightarrow} \langle \vec{X}(t), \vec{Y}(t) \rangle$ 内积保持不变.
特别, $\stackrel{(8.6)}{\Rightarrow} \vec{X}(t)$ 长度 \dots

§4.2. 测地线.

S : $\vec{r} = \vec{r}(u^1, u^2)$ 正则曲面.

$C: u^\gamma = u^\gamma(t)$ 为 S 上的光滑曲线.

定义: 若 C 的切向量场沿 C 在 S 上平行的, 则称 C 为 S 上的测地线.

$$\vec{X}(t) := \frac{d\vec{r}}{dt} = \frac{\partial \vec{r}}{\partial u^\alpha}(u^\alpha(t), u^\beta(t)) \frac{du^\alpha(t)}{dt}$$

$$\Rightarrow X^\alpha(t) = \frac{du^\alpha(t)}{dt}$$

$$\Rightarrow \text{测地方程: } \frac{d^2 u^\alpha}{dt^2} + \Gamma_{\beta\gamma}^\alpha \frac{du^\beta}{dt} \frac{du^\gamma}{dt} = 0.$$

Next time: 变分法!!!

作业: 求 $x^2 + y^2 + z^2 = 1$ 上的测地线.