

# Single-Layered MoS<sub>2</sub> Directly Grown on Rutile TiO<sub>2</sub>(110) for Enhanced Interfacial Charge Transfer

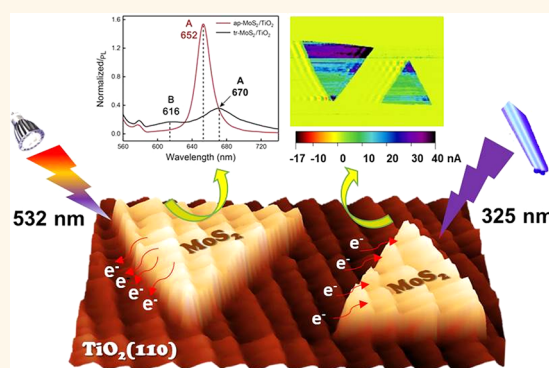
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## S Supporting Information

**ABSTRACT:** Interfacial charge transfer is critical for the photocatalytic activities of compositional photocatalysts. In this work, we have developed a strategy of growing single-layer MoS<sub>2</sub> sheets on the rutile TiO<sub>2</sub>(110) single-crystal surface using a chemical vapor deposition method. Both on-site and off-site characterizations confirmed the monolayer thickness and single crystallinity of the MoS<sub>2</sub> adlayer as well as the atomic flatness of the composite surface. Without the presence of contamination, the charge flow across the interface of MoS<sub>2</sub> and TiO<sub>2</sub> is greatly enhanced, which hence favors the charge separation under excitations and boots up the catalytic activity of the composite system. Moreover, we found the luminescing property of MoS<sub>2</sub> is significantly tailored upon coupling with the TiO<sub>2</sub> surface. Our work has established a method for revealing the interface properties of the transition-metal dichalcogenides and oxide semiconductors at the atomic level.

**KEYWORDS:** MoS<sub>2</sub>, rutile-TiO<sub>2</sub>, interface, CVD, charge transfer



As a prototypical heterostructure of transition-metal dichalcogenides (TMDs) and oxide semiconductors, the MoS<sub>2</sub>/TiO<sub>2</sub> composite has received increasing research interests due to its potential applications in electronics,<sup>1,2</sup> photovoltaics,<sup>3</sup> and in particular, photocatalysis.<sup>4,5</sup> Theoretically, the TMDs are characteristic of sufficiently high electron conductivity and mobility<sup>6</sup> while the oxides, TiO<sub>2</sub> and ZnO for instance, possess relatively large band gap yet meantime the proper anchoring sites for molecular reactions.<sup>7,8</sup> Their combination constitutes heterojunction wherein the intrinsic properties of the two materials are fairly preserved except at the interface region an dipole field is established.<sup>9–11</sup> The latter can take an active effect in separating the photoinduced carriers, which thus leads to increased visible light absorption and enhanced redox reactivities.<sup>12,13</sup>

Experimentally, the heterojunction of TMDs and oxides was usually prepared *via* hydrothermal synthesis or other wet chemical methods.<sup>14,15</sup> It is the applied solution environments that inevitably cause large amounts of adsorbates to be intercalated in the TMD/oxide interface or strongly bound on the surface, which can severely weaken the hybridization and, hence, hamper the charge transfer between TMD and oxide.

As a solution, vacuum preparation<sup>16</sup> or chemical vapor deposition (CVD) strategies<sup>17–19</sup> can perfectly avoid such problems because of their clean/inert atmosphere and high synthetic temperatures. Recently, there has much progress in preparing MoS<sub>2</sub> or analogue TMDs on various oxide surfaces *via* the CVD method.<sup>20–25</sup> However, direct synthesis of MoS<sub>2</sub> on a reducible oxide surface such as TiO<sub>2</sub> remains a challenge considering the ease of reduction and corrosion of the TiO<sub>2</sub> surface.<sup>9,10,25</sup> In this work, we report that high-quality MoS<sub>2</sub> sheets can be grown on a rutile TiO<sub>2</sub>(110) single-crystal surface by using MoO<sub>3</sub> and sulfur as the reagents of a CVD process. Rutile TiO<sub>2</sub>(110) was chosen as the substrate because of its fairly good thermal stability, well-understood surface structure, and the easy availability of single crystals, all perfectly fitting the criteria for constructing a model system. With the optimized synthetic parameters in hand, the atomic flatness of the TiO<sub>2</sub>(110) surface was well preserved, while the MoS<sub>2</sub> was fabricated with single crystallinity and monolayer thickness and

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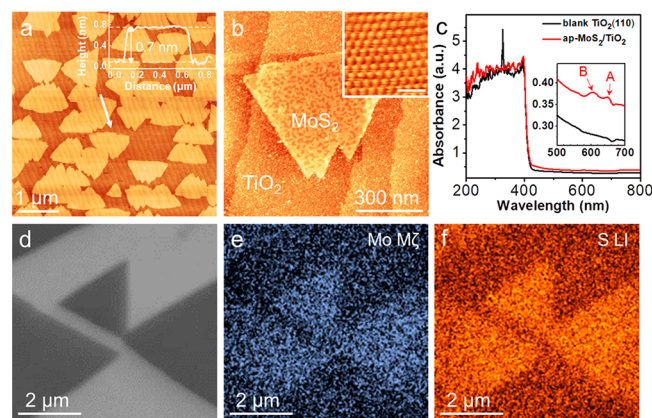
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intimately attached to the substrate. Moreover, subsequent characterizations clearly reveal that without the presence of intercalated species the interfacial charge transfer between MoS<sub>2</sub> and TiO<sub>2</sub> is substantially enhanced under various illuminating conditions, and the luminescing property of MoS<sub>2</sub> is significantly tailored. This work has deepened our understandings of the TMDs/oxide heterojunctions at an atomic level.

## RESULTS AND DISCUSSION

**Characterization of the As-Prepared MoS<sub>2</sub> Sheets on TiO<sub>2</sub>(110).** Figure 1 shows a series of characterizations of the

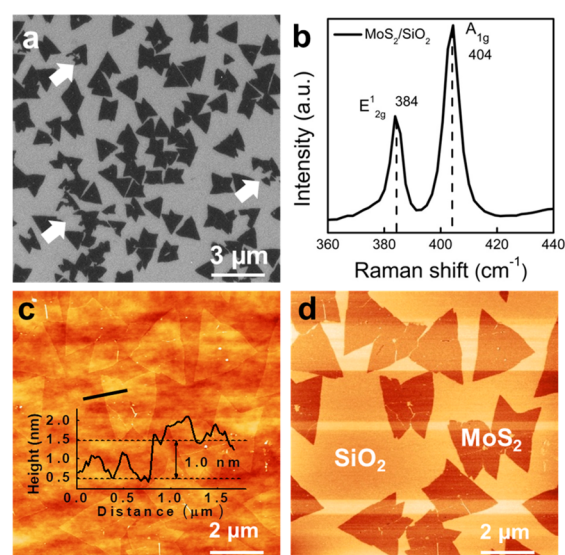


**Figure 1.** Characterizations of the as-grown MoS<sub>2</sub> sheets on an r-TiO<sub>2</sub>(110) single-crystal substrate. (a) AFM topographic image of the ap-MoS<sub>2</sub>/TiO<sub>2</sub>(110) sample surface. The inset shows the height profile of a MoS<sub>2</sub> sheet. (b) STM image of a specially small MoS<sub>2</sub> nanosheet grown on an Nb-doped r-TiO<sub>2</sub>(110) surface. The inset shows the atomic resolution. The scale bar is 1 nm. (c) UV-vis spectra of the blank TiO<sub>2</sub> (black) and the ap-MoS<sub>2</sub>/TiO<sub>2</sub> (red). Inset is the magnified spectrum of the visible-light regime where the characteristic absorption of MoS<sub>2</sub> can be identified. (d–f) SEM image and elemental mapping (e for Mo and f for S) of a selected region on the ap-MoS<sub>2</sub>/TiO<sub>2</sub> sample. The triangular dark features in (d) correspond to MoS<sub>2</sub>.

as-prepared MoS<sub>2</sub> on TiO<sub>2</sub> substrate (termed as ap-MoS<sub>2</sub>/TiO<sub>2</sub>). The atomic force microscopy (AFM) image in Figure 1a clearly shows the triangular-shaped MoS<sub>2</sub> sheets randomly dispersed on the r-TiO<sub>2</sub>(110) substrate whose atomic flatness and the step-terrace structure are well preserved. The MoS<sub>2</sub> sheets are measured as wide as 1–3 μm and as thick as ~0.7 nm, consistent with the height of a single-layer MoS<sub>2</sub>. Topographically, the synthesized MoS<sub>2</sub> is found closely attached to the TiO<sub>2</sub> surface, even at the step region. To characterize *in situ* the atomic structure of the synthesized MoS<sub>2</sub> nanosheet, we have performed the same CVD growth on a Nb-doped TiO<sub>2</sub>(110) singlecrystal substrate (shortened as Nb:TiO<sub>2</sub>(110)) and obtained identical AFM topographies as on undoped TiO<sub>2</sub> substrate, as shown in Figure S1. Figure 1b and the inset show the large-area and atomically resolved scanning tunneling microscope (STM) images of the as-prepared MoS<sub>2</sub>/Nb:TiO<sub>2</sub>(110) sample, respectively, which unambiguously confirm the atomic structure of MoS<sub>2</sub>. One may notice the corrugated pattern resolved on the MoS<sub>2</sub> patch. We tentatively attribute it to single-layer vacancies distributed at the TiO<sub>2</sub> substrate underneath the MoS<sub>2</sub> sheet since they have the same depth as a single step on TiO<sub>2</sub>(110). Moreover, Figure S2a,b clearly shows the continuous MoS<sub>2</sub> lattice

crossing both the step and the vacancy features of the TiO<sub>2</sub> substrate, demonstrating the single crystallinity of the MoS<sub>2</sub> sheet. Besides the AFM/STM measurements, the ultraviolet–visible absorption spectroscopy (UV–vis) of the ap-MoS<sub>2</sub>/TiO<sub>2</sub> sample shows slightly enhanced absorbance of visible light relative to the blank TiO<sub>2</sub>, as presented in Figure 1c. The inset shows the intentionally magnified visible light region wherein the characteristic absorptions of MoS<sub>2</sub> at around 650 and 605 nm can be clearly identified.<sup>20</sup> Figures 1d–f show the scanning electron microscopy (SEM) image as well as the corresponding elemental mapping of Mo (Figure 1e) and S (Figure 1f), which also strongly support the formation of MoS<sub>2</sub> on TiO<sub>2</sub> surface.

For better evaluating the quality of the synthesized MoS<sub>2</sub>, we performed procedures to peel off the TiO<sub>2</sub> substrate by following the reported recipe on similar TMD/oxide systems.<sup>9,26,27</sup> Figure 2a shows the SEM images of the MoS<sub>2</sub>



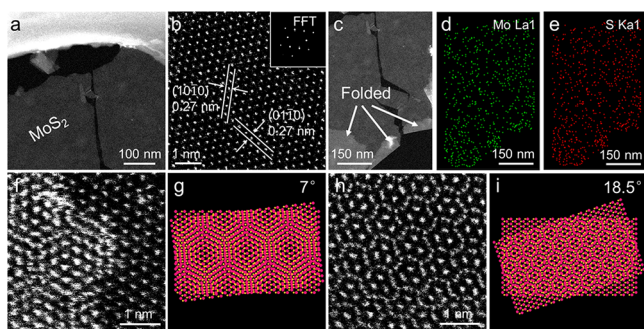
**Figure 2.** (a) SEM image and (b) Raman spectrum of the MoS<sub>2</sub> transferred onto a SiO<sub>2</sub>/Si substrate. White arrows mark the broken MoS<sub>2</sub> sheets due to the transfer operation. (c, d) AFM topography and the corresponding phase images of the MoS<sub>2</sub>/SiO<sub>2</sub>/Si surface sample. (c, inset) Height profile along the black line which runs across the edge of a MoS<sub>2</sub> sheet.

transferred onto a SiO<sub>2</sub>/Si surface, which clearly attests that the transfer operation is successful but some of the MoS<sub>2</sub> sheets (marked by the white arrows) are unavoidably broken during the transfer. Figure 2b shows the Raman spectrum taken on the transferred MoS<sub>2</sub>/SiO<sub>2</sub>/Si. The two characteristic peaks, 384 and 404 cm<sup>−1</sup>, can be clearly identified and are ascribed to the E<sub>2g</sub><sup>1</sup> and A<sub>1g</sub> phonon modes of MoS<sub>2</sub>, respectively.<sup>20,28</sup> Noticeably, the interval between E<sub>2g</sub><sup>1</sup> and A<sub>1g</sub> peaks (≤20 cm<sup>−1</sup>) strongly supports that our MoS<sub>2</sub> is perfectly single layer. We also directly measured Raman spectroscopy on the ap-MoS<sub>2</sub>/TiO<sub>2</sub> samples, but the MoS<sub>2</sub> signals are completely buried by the strong background signal of TiO<sub>2</sub>, as shown in Figure S3. Parts c and d of Figure 2 show the AFM topography and corresponding phase images of the transferred MoS<sub>2</sub> sheets on the SiO<sub>2</sub>/Si substrate. From the profile, one immediately notices that the SiO<sub>2</sub>/Si substrate possesses a much rougher surface than TiO<sub>2</sub>(110), and the transferred MoS<sub>2</sub> shows a thickness larger than 1.0 nm,



significantly higher than that of ap-MoS<sub>2</sub>/TiO<sub>2</sub>, which is possibly due to the presence of the intercalated species.

The peeled MoS<sub>2</sub> also renders the direct characterization with transmission electron microscopy (TEM). Parts a and b of Figure 3 show the large area and atomically resolved image



**Figure 3.** TEM characterization of the MoS<sub>2</sub> transferred to a copper mesh. (a, b) Typical low- and high-resolution TEM images of the prepared MoS<sub>2</sub>. (b, inset) is the FFT image showing the periodicity of the MoS<sub>2</sub> lattice. (c) Selected low-magnification image and corresponding (d) Mo and (e) S mapping of a MoS<sub>2</sub> film containing folded regions which is formed during the transfer process. (f, h) High-resolution images of different folded areas with distinct moiré patterns; corresponding models are given in (g) and (i), respectively.

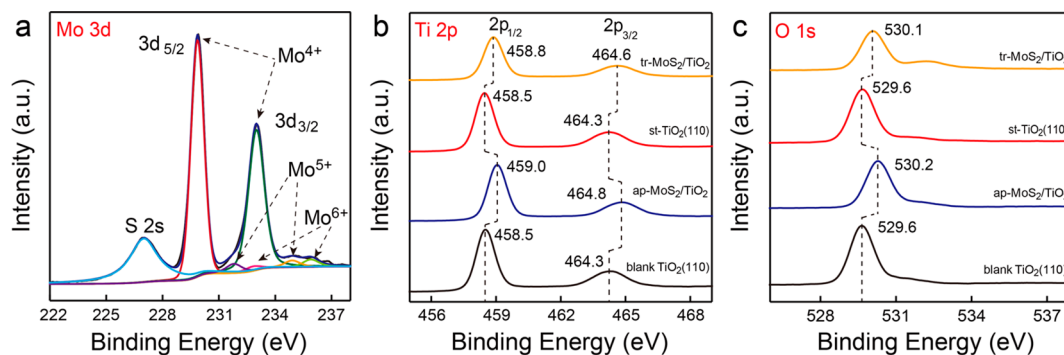
of the MoS<sub>2</sub> transferred onto a copper mesh. Both the interplane spacing measurement and the fast Fourier transformation (FFT) result in a lattice constant of 0.27 nm, which perfectly coincides with the reported value of MoS<sub>2</sub>(0001).<sup>20,29</sup> Figure 3c shows a selected edge area of MoS<sub>2</sub>, where film back-folding and overlaying can be observed. The elemental mapping of both Mo and S (parts d and e, respectively, of Figure 3) perfectly reproduces the observed shape of the MoS<sub>2</sub> sheet, once again proving the composition. For the folded region, high-resolution images (Figures 3f,h) show different moiré patterns, which can be modeled by rotational stacking of two MoS<sub>2</sub> layers with angles of 7° (Figure 3g) and 18.5° (Figure 3i), respectively.

**XPS Measurements of the Various MoS<sub>2</sub>/TiO<sub>2</sub> Samples.** Once the single-layer thickness and single crystallinity of the fabricated MoS<sub>2</sub> have been confirmed, we can examine the status of the MoS<sub>2</sub>/TiO<sub>2</sub> interface. Figure 4 shows the X-ray photoelectron spectroscopy (XPS) spectra taken on a typical ap-MoS<sub>2</sub>/TiO<sub>2</sub> sample. The full-range survey spectrum is presented in Figure S4, from which the cleanness

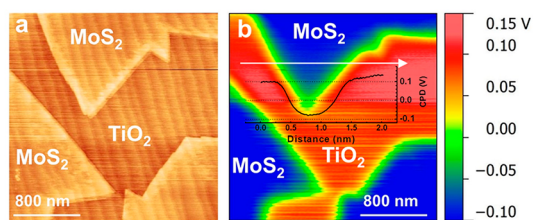
of the compositional system can be evidenced. The high-resolution Mo 3d spectrum in Figure 4a demonstrates the dominance of Mo<sup>4+</sup>, whose 3d<sub>5/2</sub> and 3d<sub>3/2</sub> were found at 229.8 and 233.0 eV, respectively, consistent with that reported for MoS<sub>2</sub> in the literature.<sup>30</sup> At the same time, trace amounts of Mo<sup>5+</sup> and Mo<sup>6+</sup> were also observed, possibly due to the unreacted MoO<sub>3</sub> particles (Figure S5) or air-induced oxidation of MoS<sub>2</sub> at the defect sites. On the lower binding energy (BE) side of Mo 3d a prominent S 2s transition is also clearly recognized, demonstrating the S component of MoS<sub>2</sub>. After MoS<sub>2</sub> was peeled off, no detectable Mo and S signals were found on the stripped TiO<sub>2</sub>(110) (termed as st-TiO<sub>2</sub>(110)) sample, as shown in Figure S6.

In Figure 4b,c we show the sequence of XPS spectra of Ti 2p and O 1s obtained on four different samples, *i.e.*, the blank TiO<sub>2</sub>(110), the ap-MoS<sub>2</sub>/TiO<sub>2</sub>, the st-TiO<sub>2</sub>(110), and the MoS<sub>2</sub>/TiO<sub>2</sub> composite obtained by transferring MoS<sub>2</sub> back onto a bare r-TiO<sub>2</sub>(110) surface (termed as tr-MoS<sub>2</sub>/TiO<sub>2</sub>). It can be clearly seen that in the presence of MoS<sub>2</sub> both Ti 2p and O 1s are shifted to higher BE synchronically by 0.5–0.6 eV, while those of blank TiO<sub>2</sub>(110) and st-TiO<sub>2</sub>(110) are essentially the same. In Figure S7a,b we show the deconvolution of the Ti 2p and O 1s spectra of ap-MoS<sub>2</sub>/TiO<sub>2</sub>, respectively. No identifiable signatures for the Ti<sup>3+</sup> species can be observed, although they were frequently observed in the composite of TiO<sub>2</sub> and MoS<sub>2</sub> fabricated by other methods. In Figure S7c,d we arbitrarily shift each Ti 2p (O 1s) spectrum to the same binding energy and normalize them to the same intensity, in order to check the profile changes. Noticeably, all four samples show the exact overlapping of the Ti 2p spectra, manifesting the absence of Ti<sup>3+</sup> formation during the CVD growth. In contrast, the O 1s of tr-MoS<sub>2</sub>/TiO<sub>2</sub> shows significantly increased OH features, dictating the introduced intercalation during the wet treatments. With these detailed analyses, the blue-shifts in the BE of both Ti 2p and O 1s of the ap-MoS<sub>2</sub>/TiO<sub>2</sub> samples can be attributed to the static field established at the interface, similar to the cases of graphene grown on TiO<sub>2</sub> and ZnO surfaces.<sup>9</sup>

The static electric field at the MoS<sub>2</sub>/TiO<sub>2</sub> interface should originate from the work function difference between MoS<sub>2</sub> and TiO<sub>2</sub>. A direct proof of such proposition comes from the Kelvin probe force microscopy (KPFM) measurements of the MoS<sub>2</sub>/TiO<sub>2</sub> sample. As shown in Figure 5, the surface potential image clearly evidences a lower contact potential ( $\sim -0.1$  V) in the MoS<sub>2</sub> region and a higher contact potential ( $\sim 0.1$  V) in the uncovered TiO<sub>2</sub> region of the ap-MoS<sub>2</sub>/TiO<sub>2</sub> sample surface, as compared with the AFM tip. This means in



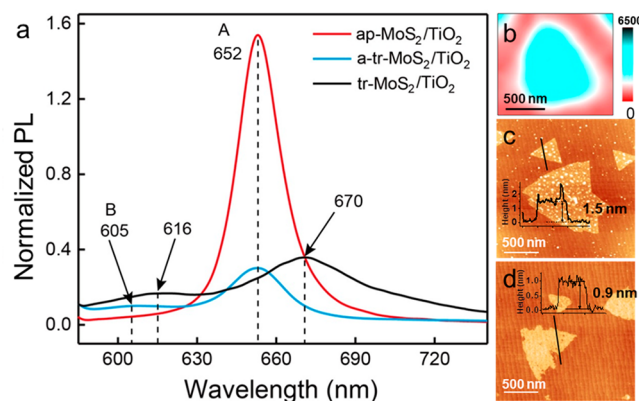
**Figure 4.** (a) High-resolution XPS spectrum of Mo 3d on the ap-MoS<sub>2</sub>/TiO<sub>2</sub> sample surface. (b, c) Series XPS spectra of Ti 2p and O 1s collected on different surfaces of blank TiO<sub>2</sub>(110), ap-MoS<sub>2</sub>/TiO<sub>2</sub>, st-TiO<sub>2</sub>(110), and tr-MoS<sub>2</sub>/TiO<sub>2</sub>, respectively.



**Figure 5.** (a) Topography and (b) corresponding surface potential image of an ap-MoS<sub>2</sub>/TiO<sub>2</sub> sample measured by KPFM. (b, inset) Line profile along the white arrow, indicating that the MoS<sub>2</sub> region has  $\sim -0.1$  V while the TiO<sub>2</sub> region has  $\sim 0.1$  V surface potential relative to the AFM tip.

the MoS<sub>2</sub>-covered region the MoS<sub>2</sub> adlayer donates electrons to the TiO<sub>2</sub> substrate across the interface due to the work function difference. The redistributed charges around the MoS<sub>2</sub>/TiO<sub>2</sub> interface thus construct a dipole field pointing along the surface normal. In this regard, introduction of any intercalation species may lead to a reduced interface dipole and, hence, a smaller shift of BE relative to the ap-MoS<sub>2</sub>/TiO<sub>2</sub> sample, which has a clean interface, as has been demonstrated by the series of XPS spectra in Figure 4. To support the above idea, we also conducted control experiments by measuring the KPFM of both ap-MoS<sub>2</sub>/TiO<sub>2</sub> and the tr-MoS<sub>2</sub>/SiO<sub>2</sub> samples with a different AFM tip. As shown in Figure S8, due to the blocked charge transfer the MoS<sub>2</sub> nanosheets on SiO<sub>2</sub> have obviously higher surface potential than on the TiO<sub>2</sub> substrate and show little difference ( $\sim 20$  mV) against the SiO<sub>2</sub> surface. In contrast, the ap-MoS<sub>2</sub>/TiO<sub>2</sub> sample still shows clear surface potential difference in MoS<sub>2</sub> and TiO<sub>2</sub> regions, respectively, in line with the proposed charge transfer between the two materials.

**Photoluminescence Property of the MoS<sub>2</sub>/TiO<sub>2</sub> Samples.** For MoS<sub>2</sub>, due to its semiconducting nature, the optical property is attractive and important. It can be significantly affected by the substrate, the surface, as well as the interface. In Figure 6a, we present the photoluminescence (PL) spectra collected on various different MoS<sub>2</sub>/TiO<sub>2</sub> samples. As can be seen, the PL spectrum of the tr-MoS<sub>2</sub>/TiO<sub>2</sub> sample shows two peaks centered at 670 and 616 nm upon excitation of the 532 nm laser. They correspond to the A and B excitonic optical transitions in the  $K/K'$  valleys of monolayer MoS<sub>2</sub>, which originate from the spin-orbit coupling induced band splitting.<sup>31,32</sup> In contrast, the spectrum of the ap-MoS<sub>2</sub>/TiO<sub>2</sub> sample shows a single peak which is highly symmetric and centered at 652 nm. This peak should also correspond to the A excitonic emission of MoS<sub>2</sub> yet with slight blue shift in energy ( $\sim 51$  meV) if compared with tr-MoS<sub>2</sub>/TiO<sub>2</sub> and the results from previous reports. A PL mapping of the ap-MoS<sub>2</sub>/TiO<sub>2</sub> taken at this peaking wavelength is shown in Figure 6b, demonstrating the luminescing uniformity of the fabricated MoS<sub>2</sub> nanosheet. It is well-known that the MoS<sub>2</sub> emission can be affected by the dielectric surroundings.<sup>33,34</sup> However, the built-in screening effect of TiO<sub>2</sub> substrate merely cannot explain the observed blue shift in the emission of the ap-MoS<sub>2</sub>/TiO<sub>2</sub>. Other contributions have to be taken into account, which leads us to the dipole field established at the MoS<sub>2</sub>/TiO<sub>2</sub> interface. Considering that the electron affinity of monolayer MoS<sub>2</sub> (4.5 eV)<sup>35</sup> is lower than that of TiO<sub>2</sub> (5.2 eV),<sup>36,37</sup> the perfect coupling at the interface of the ap-MoS<sub>2</sub>/TiO<sub>2</sub> sample will generate a perpendicular electric field at the interface, resulting in band renormalization



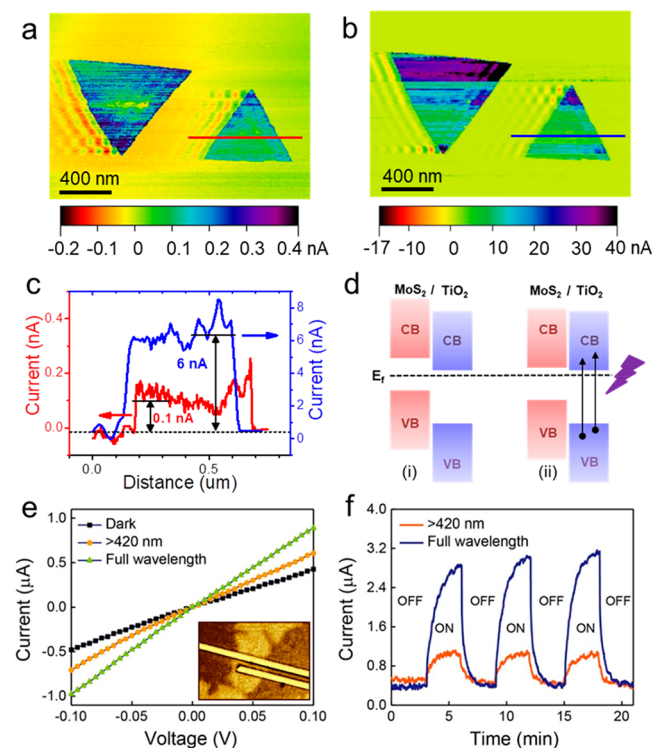
**Figure 6.** (a) Photoluminescence spectra of ap-MoS<sub>2</sub>/TiO<sub>2</sub> (red) and tr-MoS<sub>2</sub>/TiO<sub>2</sub> samples with (blue) and without (black) annealing treatment. With intercalated species present, the A and B peaks of the tr-MoS<sub>2</sub>/TiO<sub>2</sub> samples were observed to shift to lower energy. (b) Photoluminescence mapping taken with 652 nm wavelength on a MoS<sub>2</sub> sheet of the ap-MoS<sub>2</sub>/TiO<sub>2</sub> sample. (c, d) AFM images of a typical tr-MoS<sub>2</sub>/TiO<sub>2</sub>(110) sample before (c) and after (d) thermal annealing at 400 °C in nitrogen atmosphere. (c, d, insets) Profiles across individual MoS<sub>2</sub> sheets demonstrating that the annealing treatment is able to remove most of the surface adsorbates and the intercalated species but can also lead to the corruption of MoS<sub>2</sub>.

in MoS<sub>2</sub>. It has been predicted that applying an electric field perpendicular to the 2D plane can modify the on-site energy of atoms in MoS<sub>2</sub>, which tailors its electronic structure.<sup>38</sup> Under the positive perpendicular electric field, the bandgap at  $K/K'$  valleys will be enlarged, resulting in the blue shift of the PL peaks. As a control experiment, we show that the perfect coupling at the MoS<sub>2</sub>/TiO<sub>2</sub> interface can be re-established by annealing after the peeling off and transfer processes. By additionally annealing the tr-MoS<sub>2</sub>/TiO<sub>2</sub> sample (hence termed as a-tr-MoS<sub>2</sub>/TiO<sub>2</sub>), we found that the PL spectrum gradually approached the ap-MoS<sub>2</sub>/TiO<sub>2</sub> case, showing the A peak at  $\sim 652$  nm and concomitantly an observable B peak at  $\sim 605$  nm. Correspondingly, the AFM measurements shown in Figures 6c,d clearly confirm that the postannealing treatment can efficiently remove most of the surface adsorbates and the intercalated species (see the comparison of the profiles), thus re-establishing the interface between MoS<sub>2</sub> and TiO<sub>2</sub>. It also significantly degrades the MoS<sub>2</sub> sheets by creating a considerable number of defects. With the above statements, the blueshift of the A excitonic emission relative to the intrinsic MoS<sub>2</sub> confirms the existence of the dipole field in the ap-MoS<sub>2</sub>/TiO<sub>2</sub> interface. We also measured the PL of the ap-MoS<sub>2</sub>/TiO<sub>2</sub> under the excitation of UV light with 325 nm wavelength as shown in Figure S9. The spectrum shows a single peak at around 660 nm yet with much lower intensity. In this case, the TiO<sub>2</sub> substrate can be excited. The photoinduced electrons in the TiO<sub>2</sub> substrate may donate back into MoS<sub>2</sub>, which hence reduces the dipole field at the interface. As a consequence, the PL peak from MoS<sub>2</sub> greatly weakens and redshifts slightly.

**Photocurrent Responses of the MoS<sub>2</sub>/TiO<sub>2</sub> Composite.** Finally, the most expected consequence of a clean MoS<sub>2</sub>/TiO<sub>2</sub> heterojunction is the improved photoactivity of the composite. With our model system, we can directly investigate the interfacial charge-transfer behavior upon light irradiation at the absence of interface contaminations. Parts a and b of



Figures 7 show the tunneling AFM (TUNA) measurements of an ap-MoS<sub>2</sub>/TiO<sub>2</sub> sample prepared on a 0.5%Nb-doped



**Figure 7.** Tunneling AFM measurements of the MoS<sub>2</sub> sheet prepared on a 0.5% Nb-doped TiO<sub>2</sub>(110) substrate under (a) dark conditions and (b) continuous irradiation of UV light. (c) Current profiles along the red and blue lines in (a) and (b), respectively. (d) Schematic model shows the electron excitation in TiO<sub>2</sub> upon UV light irradiation and the subsequent band realignment at the interface. CB and VB denote conduction and valence bands, respectively. (e) Current–voltage curves measured with two electrodes directly deposited on a MoS<sub>2</sub> sheet. (Inset) SEM image of the setup. (f) Current responses to the switching on/off of the visible (orange) and full-range (blue) light irradiations. The bias was fixed at 100 mV.

TiO<sub>2</sub>(110) substrate under dark and UV light illumination conditions, respectively. From the profiles shown in Figure 7c one can clearly see that even under dark conditions the MoS<sub>2</sub>-covered area already shows increased out-of-plane conductivity compared to the bare Nb:TiO<sub>2</sub> surface. Upon switching on the UV light with a power of 5 W and centered wavelength of 365 nm, the conductivity of the MoS<sub>2</sub> area is dramatically enhanced by almost 2 orders of magnitude, while the uncovered TiO<sub>2</sub> remains more or less the same as in darkness. In a control experiment shown in Figure S10, another MoS<sub>2</sub> sheet is imaged with and without a normal illuminating torch (visible light) switched on, which shows an ignorable difference. Therefore, the observed significant increase of the tunneling current under UV light can be ascribed to the sufficient charge transfer from TiO<sub>2</sub> into MoS<sub>2</sub> since TiO<sub>2</sub> is a bulk material and can be readily excited under UV irradiations. In this case, the energy difference between the conduction bands of MoS<sub>2</sub> and TiO<sub>2</sub> may be reduced, which thus facilitates the charge transfer from TiO<sub>2</sub> to MoS<sub>2</sub>, as presented by the schematic model in Figure 7d. Better understanding of the MoS<sub>2</sub>/TiO<sub>2</sub> heterojunction can be obtained from the in-plane conductivity changes under different irradiation conditions. As shown in

Figure 7e and the inset, two gold electrodes with interval of 1  $\mu\text{m}$  were deposited onto a single MoS<sub>2</sub> sheet, and the current responses were collected during ramping of the bias. In this case, the irradiations were provided by a xenon lamp with power of 150 W. The conductance of the MoS<sub>2</sub>/TiO<sub>2</sub> composite was measured to be 4.4, 6.4, and 9.3 S·m<sup>-1</sup> for dark, visible-light, and full-length irradiations, respectively. Figure 7f shows the current signal changes upon switching the light on/off repeatedly, which also presents an on/off response except for the relatively slow response speed. The related charge-transfer dynamics will be investigated with time-resolved photoelectron spectroscopy and reported in a separated work. Nevertheless, the above measurements clearly demonstrate that the photoexcited electrons in TiO<sub>2</sub> are efficiently transferred into MoS<sub>2</sub> without interruption of the intercalated species. Consequently, the separated photoelectrons and holes, respectively, render the reduction and oxidation reactions at the corresponding surface sites.

## CONCLUSION

In conclusion, by using MoO<sub>3</sub> and sulfur powder as solid precursors for CVD growth, we have successfully synthesized high-quality single-layer MoS<sub>2</sub> sheets on the atomically flat r-TiO<sub>2</sub>(110) single-crystal surface. Detailed characterizations including AFM/STM, XPS, and PL measurements not only confirmed the seamless contact between MoS<sub>2</sub> and TiO<sub>2</sub> but also unambiguously revealed the efficient electron transfer across the intimate MoS<sub>2</sub>/TiO<sub>2</sub> heterojunction under different irradiation conditions. Our work has provided a strategy of establishing proper TMD/oxide model systems for a deepened understanding of the corresponding photocatalytic mechanisms.

## MATERIALS AND METHODS

**Materials.** (1) Rutile TiO<sub>2</sub>(110) (r-TiO<sub>2</sub>(110), 10 × 5 × 0.5 mm<sup>3</sup>, 99.99%, Hefei Kejing Material Technology Co., Ltd.). (2) 0.5% Nb-doped r-TiO<sub>2</sub>(110) (10 × 3 × 0.5 mm<sup>3</sup>, Shinkosha Co., Ltd.). (3) Silica substrate (SiO<sub>2</sub>/Si, Hefei Kejing Material Technology Co., Ltd.). (4) MoO<sub>3</sub> (Aladdin, 99.6%). (5) Sulfur powder (Aladdin, 99.6%). (6) N<sub>2</sub> (Linde industrial gas, 99.999%). (7) Hydrofluoric acid (HF, Sinopharm Chemical Reagent Co., Ltd.). (8) Acetone (Sinopharm Chemical Reagent Co., Ltd., 99.5%). (9) Deionized water (Millipore, Conductivity 18.2 M $\Omega$ ·cm). All of the chemicals mentioned above were used as received.

**Preparation of the Atomically Smooth r-TiO<sub>2</sub> (110) Surface.** Atomically smooth surfaces of r-TiO<sub>2</sub>(110) were routinely achieved by following a procedure reported by Yamamoto *et al.*,<sup>39</sup> which proceeds sequentially through ultrasonic cleaning with acetone and ultrapure water, etching with 20% HF for 10 min, rinsing with water, drying in a nitrogen stream, and annealing at 800 °C for 3 h in air. The atomic flatness and cleanness of the prepared surfaces were confirmed by AFM and XPS measurements repeatedly. All chemical reagents were directly used without further purification.

**CVD Growth of MoS<sub>2</sub>.** The CVD growth of MoS<sub>2</sub> on the clean r-TiO<sub>2</sub>(110) single crystal was performed in a tube furnace (OTF-1200X, Hefei Kejing Material Technology Co., Ltd.) with a dual temperature zone. The specific steps are as follows: (1) Place 120 mg of sulfur powder and 6 mg of MoO<sub>3</sub> in a separated quartz boat and place them, respectively, at the low temperature (set to 200 °C) and high temperature zone (set to 700 °C) of the tube furnace. (2) Put the r-TiO<sub>2</sub>(110) substrate upside down in the same quartz boat approximately 2 cm before the MoO<sub>3</sub> powder. (3) Flush the furnace tube with N<sub>2</sub> gas (flow rate ~300 ccm) for 20 min to completely remove the residual oxygen air. (4) Heat the MoO<sub>3</sub> zone to 100 °C under N<sub>2</sub> flow for 1 h to remove the adsorbed water. (4) Reduce the flow rate of N<sub>2</sub> to 10 sccm and heat the MoO<sub>3</sub> zone with a heating

rate of 15 °C/min. (5) When the MoO<sub>3</sub> zone reaches ~380 °C, the temperature of the sulfur zone begins to rise slowly. Finally the MoO<sub>3</sub> zone reaches 700 °C, and the sulfur zone reaches 200 °C. (6) Keep both zones at the target temperature for 10 min before turning off the heating. (7) Let the furnace cool down naturally until it reaches approximately 580 °C. Afterward, increase the N<sub>2</sub> flow rate to 300 ccm and open the shroud of the furnace to speed the cooling to room temperature. With this procedure, single-layered MoS<sub>2</sub> with a crystal size of 1–5 μm can be prepared.

**Characterization Methods.** All atomic force microscopy measurements were performed on a Multimode AFM V produced by Bruker under ambient conditions. The images were obtained with commercial silicon cantilevers from the same company. In particular, topographic images were measured with a standard silicon cantilever. Kelvin probe force microscopy (KPFM) was measured with As-doped silicon tips with different coatings (SCM-PIT-V2 or MESP, Bruker Co.). Tunneling AFM (TUNA) was measured with an As-doped silicon tip with resistivity of 0.01–0.025 Ohm-cm (SCM-PIC, Veeco Co.). For TUNA measurements under illumination conditions, electric torches with similar power (~5 W) but different center wavelengths were applied to mimic the visible-light and UV-light irradiations, respectively.

Scanning tunneling microscopy (STM) was measured with a commercial low-temperature STM (Createc) operated at liquid nitrogen temperature. The samples were slightly degassed to around 100 °C in the fast-entry chamber with pressure of  $1 \times 10^{-7}$  mbar before being transferred into STM. The UV–vis absorption spectra were acquired on a SolidSpec-3700 spectrometer (Shimadzu). TEM images were obtained by ARM-200F (JEOL). SEM measurements were performed using a SIRION 200 (FEI) model instrument. The XPS was measured with an AXIS Ultra (Kratos Analytical) instrument equipped with an Al K $\alpha$  source (1486.7 eV). PL measurements were conducted using an Iso Plan SCT 320 (Princeton) spectrometer. An excitation laser of 532 nm for the wavelength and 1.5 mW for the power was applied. The laser spot has a diameter as small as 1 μm. Raman spectroscopy was measured using a LabRamHR Raman spectrometer with an excitation light of 514.5 nm.

## ASSOCIATED CONTENT

### Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acsnano.9b02608.

Additional characterization results of the MoS<sub>2</sub>/TiO<sub>2</sub> samples including AFM, XPS, PL, and TUNA measurements (PDF)

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### Notes

The authors declare no competing financial interest.

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