

Outlines of Quantum Physics

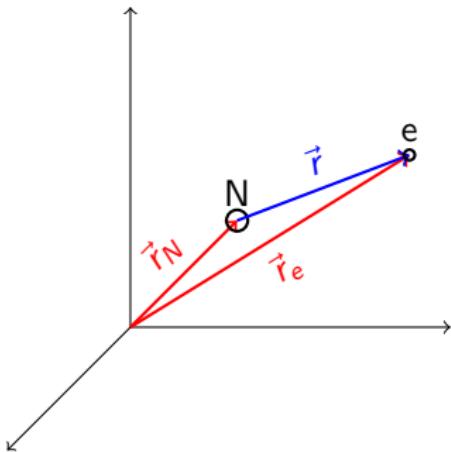
① Wave-Particle Duality

② The Schrödinger Equation

③ The Hydrogen Atom

- Schrödinger Eq. of the Hydrogen Atom
- Noninteracting Particles and Separation of Variables
- The One-Particle Central Force Problem
- Energy levels of the Hydrogenlike atom
- Hydrogenlike-Atom Wave Functions

Schrödinger Eq. of the Hydrogen Atom



氢原子或类氢原子：

$$H = -\frac{\hbar^2}{2M_N} \nabla_N^2 - \frac{\hbar^2}{2m_e} \nabla_e^2 - \frac{Ze'^2}{r}$$

其中 $\vec{r} = \vec{r}_e - \vec{r}_N$

$$e'^2 = e^2 / 4\pi\epsilon_0$$

Noninteracting Particles and Separation of Variables

$$\hat{H} = \hat{H}_1 + \hat{H}_2, \quad (\hat{H}_1 + \hat{H}_2)\Psi(q_1, q_2) = E\Psi(q_1, q_2)$$

其中 \hat{H}_1 只和粒子 1 有关; \hat{H}_2 只和粒子 2 有关;

分离变量法, 假设 $\Psi(q_1, q_2) = G_1(q_1)G_2(q_2)$:

$$\hat{H}_1 G_1(q_1)G_2(q_2) + \hat{H}_2 G_1(q_1)G_2(q_2) = EG_1(q_1)G_2(q_2)$$

$$G_2(q_2)\hat{H}_1 G_1(q_1) + G_1(q_1)\hat{H}_2 G_2(q_2) = EG_1(q_1)G_2(q_2)$$

$$\frac{\hat{H}_1 G_1(q_1)}{G_1(q_1)} + \frac{\hat{H}_2 G_2(q_2)}{G_2(q_2)} = E$$

Left part must be a constant,

$$\hat{H}_1 G_1(q_1) = E_1 G_1(q_1), \quad \hat{H}_2 G_2(q_2) = E_2 G_2(q_2)$$

$$E_1 + E_2 = E$$

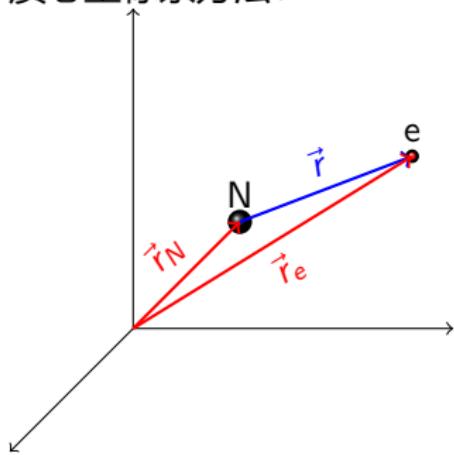
Noninteracting Particles and Separation of Variables

可推广到 n 个独立“粒子”问题：

$$\begin{aligned}\hat{H} &= \hat{H}_1 + \hat{H}_2 + \cdots + \hat{H}_n \\ \Psi(q_1, q_2, \dots, q_n) &= G_1(q_1)G_2(q_2)\cdots G_n(q_n) \\ E &= E_1 + E_2 + \cdots + E_n \\ \hat{H}_i G_i &= E_i G_i, \quad i = 1, 2, \dots, n\end{aligned}$$

Reduction of the Two-Particle Problem to Two One-Particle Problem

质心坐标系方法:



$$\begin{aligned}\vec{R}(X, Y, Z) &= \frac{m_1 \vec{r}_1 + m_2 \vec{r}_2}{m_1 + m_2} \\ \vec{r}(x, y, z) &= \vec{r}_1 - \vec{r}_2 \\ M = m_1 + m_2; \quad \mu &= \frac{m_1 m_2}{m_1 + m_2}\end{aligned}$$

$$\begin{aligned}\nabla_R^2 &= \frac{\partial^2}{\partial X^2} + \frac{\partial^2}{\partial Y^2} + \frac{\partial^2}{\partial Z^2} \\ \nabla_r^2 &= \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}\end{aligned}$$

$$\Rightarrow \frac{1}{m_1} \nabla_1^2 + \frac{1}{m_2} \nabla_2^2 = \frac{1}{M} \nabla_R^2 + \frac{1}{\mu} \nabla_r^2$$

则 S.eq 可化为：

$$\underline{\left[-\frac{\hbar^2}{2M} \nabla_R^2 - \frac{\hbar^2}{2\mu} \nabla_r^2 + V(r) \right] \Psi} = E_{total} \Psi$$

分离变量法，

$$\begin{aligned}\Psi(\vec{R}, \vec{r}) &= \Psi_c(\vec{R}) \Psi_e(\vec{r}) \\ -\frac{\hbar^2}{2M} \nabla_R^2 \Psi_c(\vec{R}) &= E_c \Psi_c(\vec{R}) \\ \left[-\frac{\hbar^2}{2\mu} \nabla_r^2 + V(r) \right] \Psi_e(\vec{r}) &= E \Psi_e(\vec{r}) \\ E_{total} &= E_c + E\end{aligned}$$

即分解成质心的平动和内部运动。
这个方法是精确的，但无法推广到多体问题。

三维中心力场

$$\left[-\frac{\hbar^2}{2\mu} \nabla^2 + V(r) \right] \Psi(\vec{r}) = E \Psi(\vec{r})$$

$$\nabla^2 = \frac{\partial^2}{\partial r^2} + \frac{2}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \left[\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right],$$

$$\hat{L}^2 = -\hbar^2 \left[\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right],$$

$$\hat{L}_z = -i\hbar \frac{\partial}{\partial \phi},$$

$$\nabla^2 = \frac{1}{r} \frac{\partial^2}{\partial r^2} r - \frac{1}{r^2 \hbar^2} \hat{L}^2$$

$$\hat{H} = -\frac{\hbar^2}{2\mu r} \frac{\partial^2}{\partial r^2} r + \frac{1}{2\mu r^2} \hat{L}^2 + V(r)$$

分离变量法求解

$$\left[-\frac{\hbar^2}{2\mu r} \frac{\partial^2}{\partial r^2} r + V(r) + \frac{1}{2\mu r^2} \hat{L}^2 \right] \Psi(r, \theta, \phi) = E \Psi(r, \theta, \phi)$$

可以分解为：

$$\Psi = R(r) Y(\theta, \phi)$$

角向方程

$$\hat{L}^2 Y(\theta, \phi) = \beta \hbar^2 Y(\theta, \phi)$$

径向方程

$$\left[\frac{1}{r} \frac{d^2}{dr^2} r + 2 \frac{\mu}{\hbar^2} (E - V) - \frac{\beta}{r^2} \right] R_\ell(r) = 0$$

角向方程，球谐函数

$$\hat{L}^2 Y(\theta, \phi) = \beta \hbar^2 Y(\theta, \phi)$$

其解

$$Y_{\ell m}(\theta, \phi) = (-1)^{\frac{m+|m|}{2}} \sqrt{\frac{(2\ell+1)(\ell-|m|)!}{4\pi(\ell+|m|)!}} P_{\ell}^{|m|}(\cos\theta) \frac{1}{\sqrt{2\pi}} e^{im\phi}$$

$$\hat{L}^2 Y_{\ell m}(\theta, \phi) = \ell(\ell+1)\hbar^2 Y_{\ell m}(\theta, \phi)$$

$$\int_0^{2\pi} d\phi \int_0^\pi d\theta Y_{\ell m}^* Y_{\ell' m'} \sin\theta = \delta_{\ell\ell'} \delta_{mm'}$$

Radial Function and Energy levels

$$\frac{\hbar^2}{2\mu}(R_\ell'' + \frac{2}{r}R_\ell') + V(r)R_\ell + \frac{\ell(\ell+1)\hbar^2}{2\mu r^2}R_\ell = ER_\ell(r)$$

一般中心力场中，其解：

$$R = R_{n\ell}(r), \quad \int_0^\infty R_{n\ell} R_{n'\ell} r^2 dr = \delta_{nn'} \\ E = E_{n\ell}$$

能量本征态一般是 $2\ell + 1$ 重简并，与 m 量子数无关。
对于类氢离子 $V(r) = Ze'^2/r$ ，其解

$$R_{n\ell}(r) = N_{n\ell} e^{-\frac{Zr}{na}} \left(\frac{2Z}{na}r\right)^\ell L_{n+\ell}^{2\ell+1} \left(\frac{2Z}{na}r\right) \\ a \equiv \hbar^2/\mu e'^2 = \frac{\hbar}{\alpha\mu c}$$

$$e'^2 = e^2 / 4\pi\epsilon_0$$

Energy levels of hydrogenlike atom

$$\begin{aligned}
 \Psi_{n\ell m} &= R_{n\ell}(r) Y_{\ell m}(\theta, \phi) \\
 H\Psi_{n\ell m} &= E_n \Psi_{n\ell m} \\
 E_n &= -\frac{Z^2}{2n^2} \cdot \frac{\mu e'^4}{\hbar^2} = -\frac{Z^2}{2n^2} \cdot \mu c^2 \alpha^2 \\
 L^2 \Psi_{n\ell m} &= \ell(\ell+1)\hbar^2 \Psi_{n\ell m} \\
 L_z \Psi_{n\ell m} &= m_\ell \hbar \Psi_{n\ell m} \\
 \ell &= 0, 1, 2, 3, \dots, n-1 \\
 m &= 0, \pm 1, \pm 2, \dots, \pm \ell
 \end{aligned}$$

- $\ell = 0$ (sharp), 1(principle), 2(diffuse), 3(fundamental), ...
- $hcR_\infty \equiv \frac{1}{2}\alpha^2 m_e c^2 \approx 13.6\text{eV}$, $a_0 \equiv \frac{\hbar}{\alpha m_e c} \approx 0.529\text{\AA}$, $\alpha \equiv \frac{(e^2)}{4\pi\epsilon_0\hbar c} \approx \frac{1}{137}$
- $R_\infty = 109\ 737.315\ 681\ 60(21) \text{ cm}^{-1}$, $R_M = R_\infty \frac{M}{m_e + M}$

Hydrogenlike-Atom Wave Functions

$$\Psi_{n\ell m} = R_{n\ell}(r) Y_{\ell m}(\theta, \phi)$$

$$R_{n\ell}(r) = \sqrt{\left(\frac{2Z}{na}\right)^3 \frac{(n-\ell-1)!}{2n[(n+\ell)!]^3}} e^{-\frac{Zr}{na}} \left(\frac{2Z}{na}r\right)^\ell L_{n+\ell}^{2\ell+1} \left(\frac{2Z}{na}r\right)$$

$$Y_{\ell m}(\theta, \phi) = (-1)^{\frac{m+|m|}{2}} \sqrt{\frac{(2\ell+1)(\ell-|m|)!}{4\pi(\ell+|m|)!}} P_\ell^{|m|}(\cos \theta) \frac{1}{\sqrt{2\pi}} e^{im\phi}$$

where $a \equiv 4\pi\epsilon_0\hbar^2/\mu e^2$

Radial Hydrogenic Wavefunction $R_{n\ell}$

$$R_{1s} = \left(\frac{Z}{a}\right)^{3/2} 2e^{-\rho}$$

$$R_{2s} = \left(\frac{Z}{2a}\right)^{3/2} 2(1 - \rho)e^{-\rho}$$

$$R_{2p} = \left(\frac{Z}{2a}\right)^{3/2} \frac{2}{\sqrt{3}} \rho e^{-\rho}$$

$$R_{3s} = \left(\frac{Z}{3a}\right)^{3/2} 2\left(1 - 2\rho + \frac{2}{3}\rho^2\right) e^{-\rho}$$

$$R_{3p} = \left(\frac{Z}{3a}\right)^{3/2} \frac{4\sqrt{2}}{3} \rho \left(1 - \rho/2\right) e^{-\rho}$$

$$R_{3d} = \left(\frac{Z}{3a}\right)^{3/2} \frac{2\sqrt{2}}{3\sqrt{5}} \rho^2 e^{-\rho}$$

$$a \equiv 4\pi\epsilon_0\hbar^2/\mu e^2$$

$$\rho = \frac{Zr}{na}$$

$\ell = 0$ 时, 在 $r = 0$ 原点处:

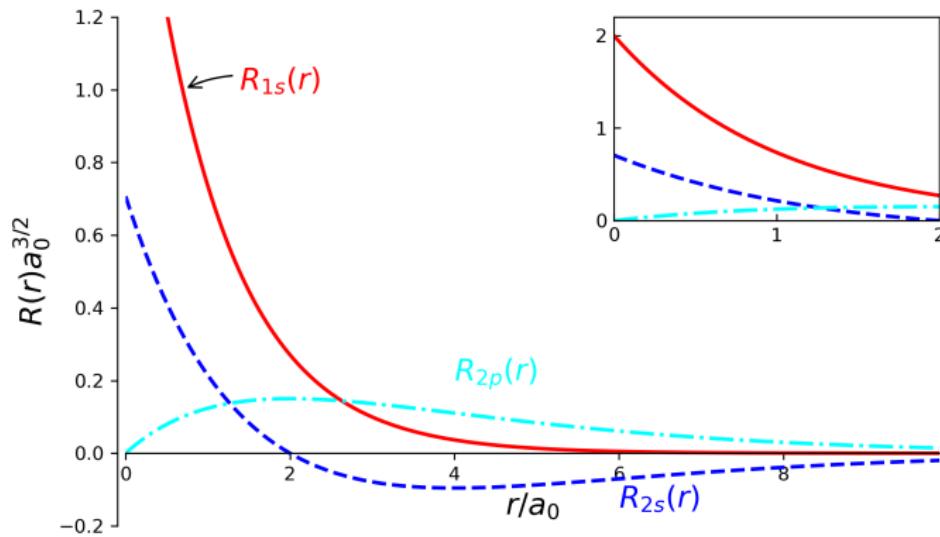
$$|\Psi_{n,\ell=0}(0)|^2 = \frac{1}{\pi} \left(\frac{Z}{na}\right)^3$$

$$\left\langle \frac{1}{r} \right\rangle = \int_0^\infty \frac{1}{r} R_{n\ell}^2 r^2 dr$$

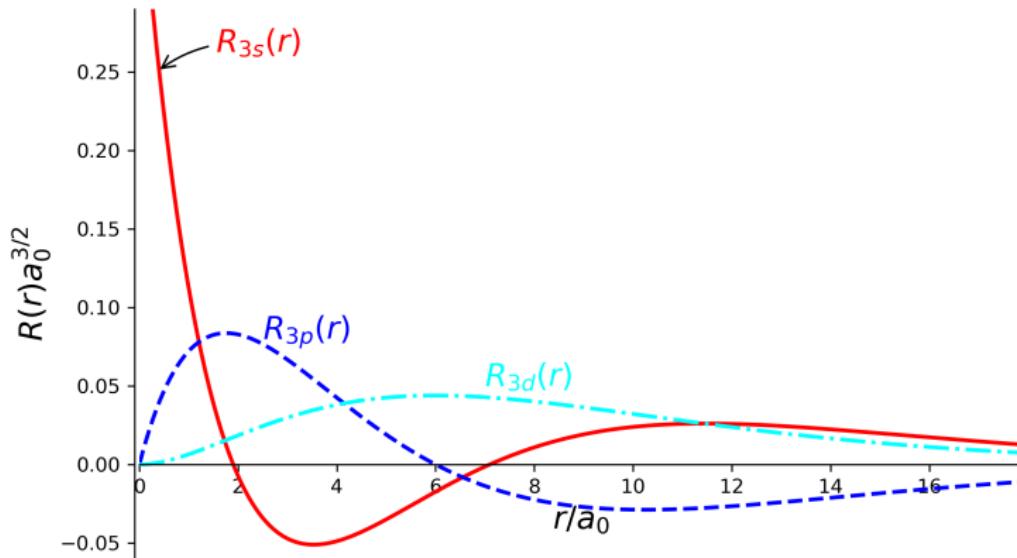
$$= \frac{1}{n^2} \left(\frac{Z}{a}\right)$$

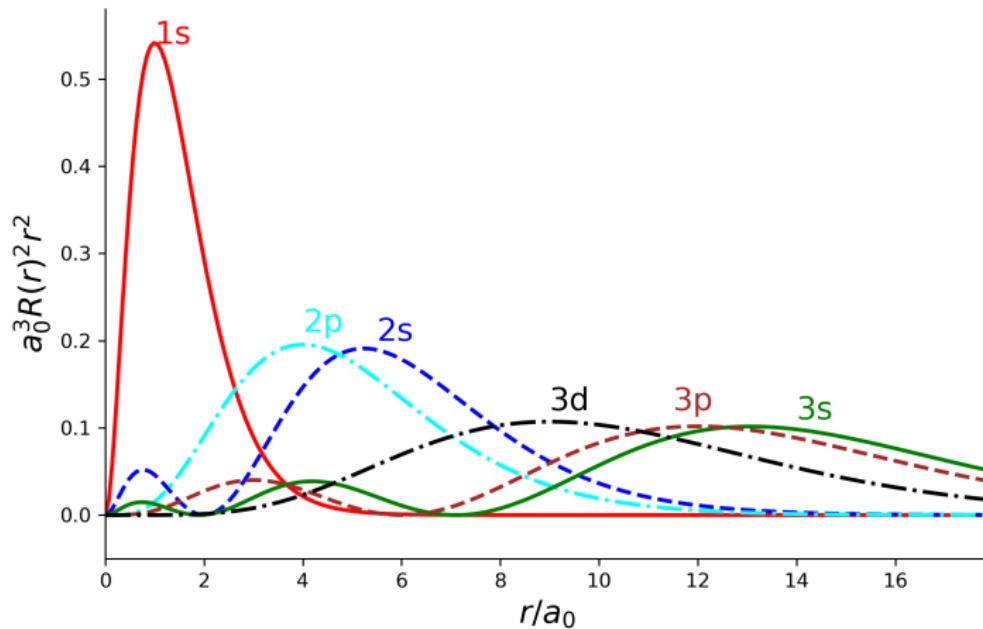
$$\left\langle \frac{1}{r^3} \right\rangle = \frac{1}{\ell(\ell+1/2)(\ell+1)} \left(\frac{Z}{na}\right)^3$$

R_{1s} & R_{2s} & R_{2p}



Electron in the Nucleus?!

R_{3s} & R_{3p} & R_{3d} 

Radial Distribution Function $R_{n\ell}^2 r^2$ 

Real Hydrogenlike Functions

$$\Psi_{n,\ell,m} = R_{n\ell}(r) S_{\ell m}(\theta) \frac{1}{\sqrt{2\pi}} e^{im\phi}$$

combinations of $\Psi_{n,\ell,m}$ and $\Psi_{n,\ell,-m}$:

$$\Psi_{n,\ell,|m|} = R_{n\ell}(r) S_{\ell,|m|}(\theta) \frac{1}{\sqrt{2\pi}} \begin{cases} \cos |m|\phi \\ \sin |m|\phi \end{cases}$$

$$\Psi_{2p_z} = \Psi_{2p_0} = \frac{1}{\sqrt{\pi}} \left(\frac{Z}{2a} \right)^{5/2} e^{-Zr/2a} z$$

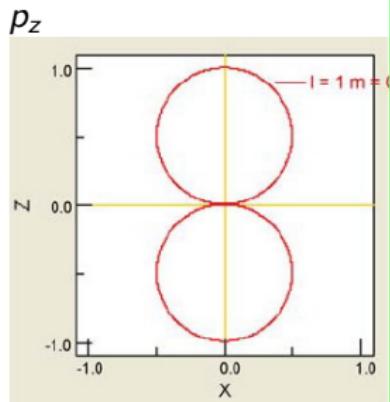
$$\Psi_{2p_x} = \frac{1}{\sqrt{2}} (\Psi_{2p_1} + \Psi_{2p_{-1}}) = \frac{1}{\sqrt{\pi}} \left(\frac{Z}{2a} \right)^{5/2} e^{-Zr/2a} x$$

$$\Psi_{2p_y} = \frac{1}{i\sqrt{2}} (\Psi_{2p_1} - \Psi_{2p_{-1}}) = \frac{1}{\sqrt{\pi}} \left(\frac{Z}{2a} \right)^{5/2} e^{-Zr/2a} y$$

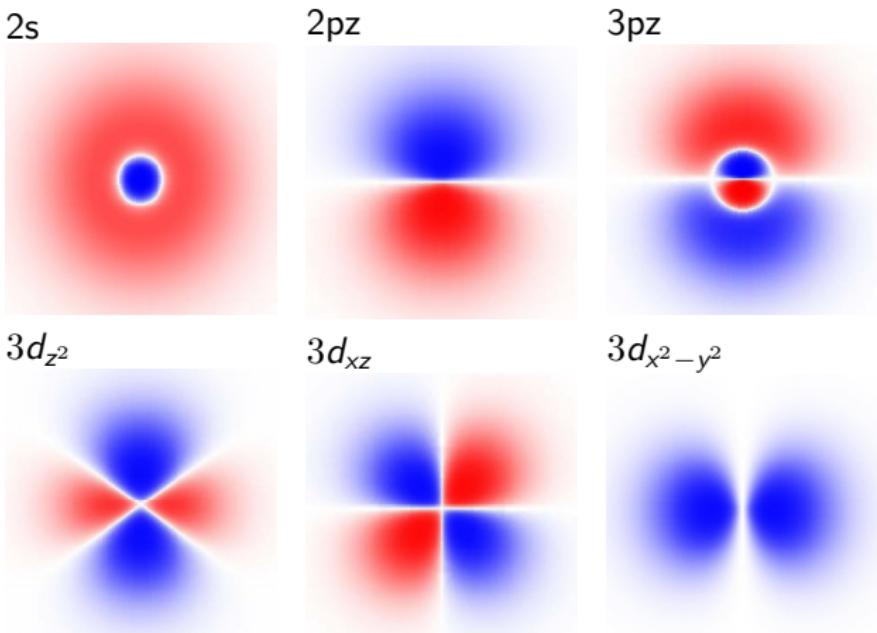
$$\Psi_{3d_{xy}} = \frac{1}{i\sqrt{2}} (\Psi_{3d_2} - \Psi_{3d_{-2}}) = \frac{2}{81\sqrt{2\pi}} \left(\frac{Z}{a} \right)^{7/2} e^{-Zr/3a} xy$$

Graphs of Real Hydrogenlike Orbitals

- Angular wave function, $|Y_{\ell m}(\theta, \phi)|$



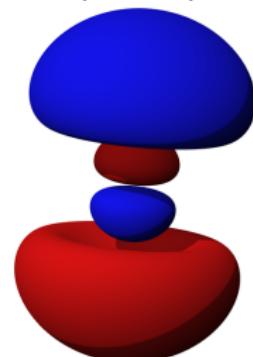
- Probability density, $|\Psi|^2$, $|\Psi| (@y = 0)$



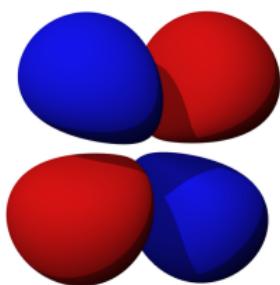
A Java viewer from: <http://www.falstad.com/qmatom/>

Graphs of Real Hydrogenlike Orbitals

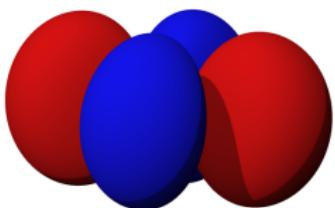
$3p_z$ ($m = 0$)



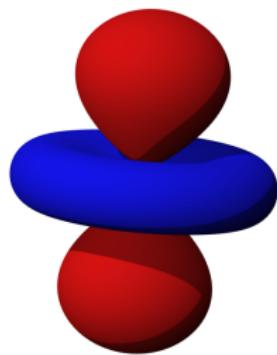
$3d_{xz}$ ($m = \pm 1$)



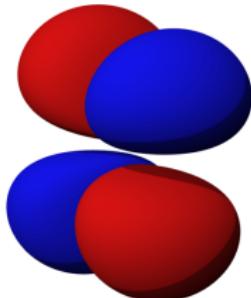
$3d_{xy}$ ($m = \pm 2$)



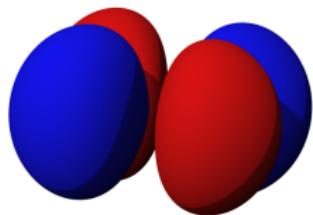
$3d_{z^2}$ ($m = 0$)



$3d_{yz}$ ($m = \pm 1$)

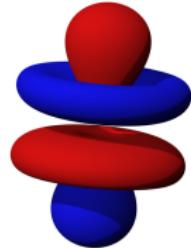


$3d_{x^2-y^2}$ ($m = \pm 2$)

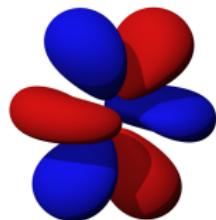


Graphs of Real Hydrogenlike Orbitals

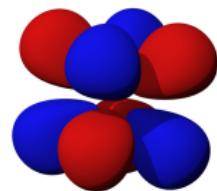
$4f_{z^3}$
($m = 0$)



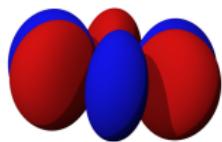
$4f_{xz^2}$
($m = \pm 1$)



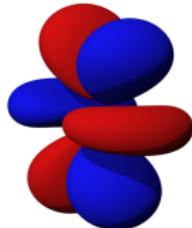
$4f_{xyz}$
($m = \pm 2$)



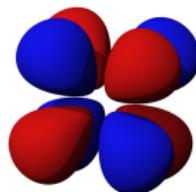
$4f_{x(x^2-3y^2)}$
($m = \pm 3$)



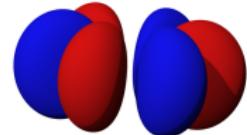
$4f_{yz^2}$
($m = \pm 1$)



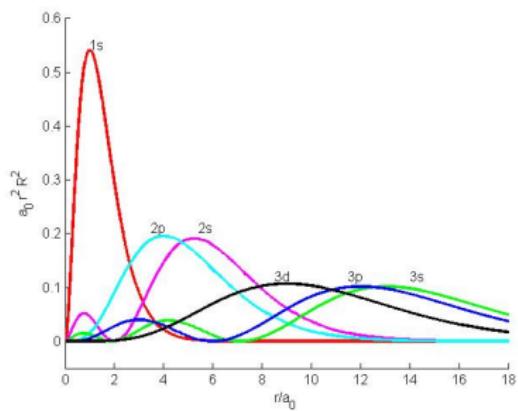
$4f_{z(x^2-y^2)}$
($m = \pm 2$)



$4f_{y(3x^2-y^2)}$
($m = \pm 3$)



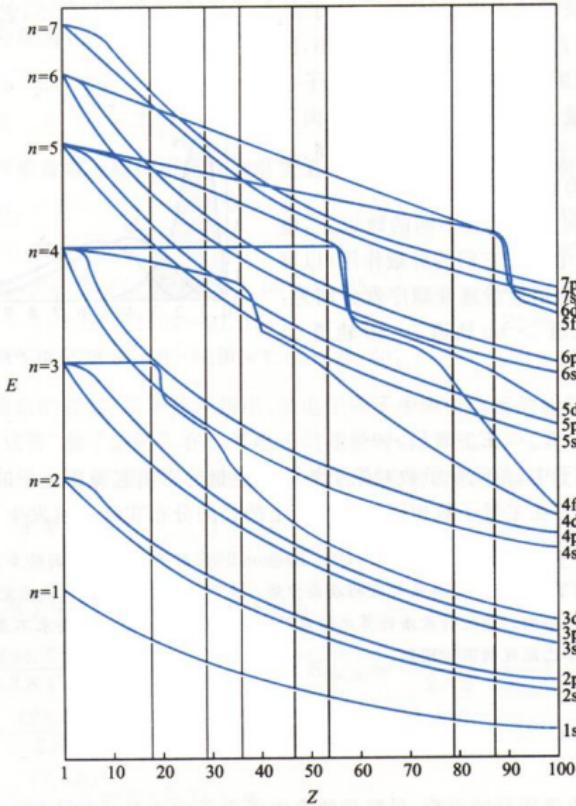
多电子原子的电子轨道能级



钻穿效应 Penetrating Effect

外层电子有穿过内层电子的布居;
 $E(ns) < E(np) < E(nd) < E(nf)$;
 能级交错。

原子轨道能级图



参考内容提纲

内容	参考书
氢原子的薛定谔方程	【Lv】§6.1 【杨】§17
多个独立粒子体系	【Lv】§6.2-6.3
单粒子三维中心力场	【Lv】§6.1 【曾】§6.1
类氢原子能级	【Lv】§6.5 【杨】§17
类氢原子波函数	【Lv】§6.6 【杨】§17