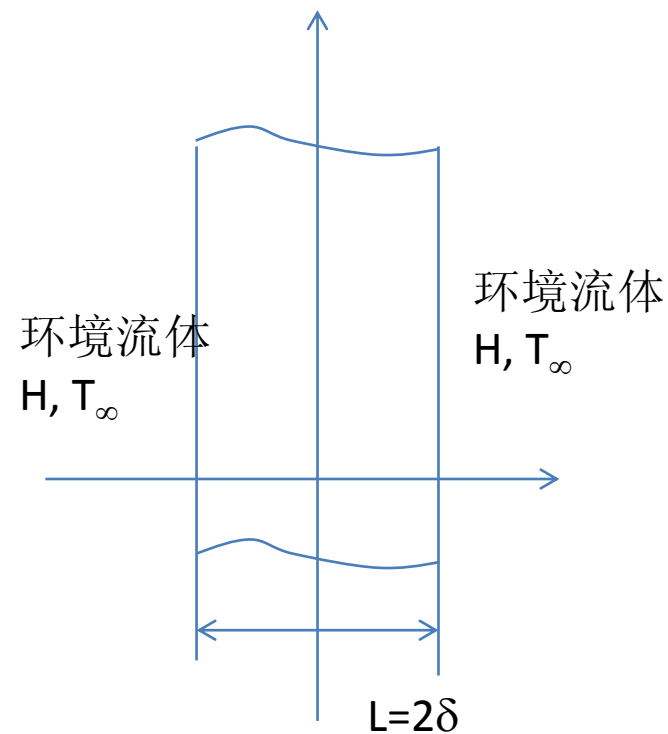


2014/3/14

内容提要

- 非稳态导热概要
- 非稳态导热的集中参数法
- 一维平板非稳态导热的分析解

3.3 一维平板非稳态导热的分析解



$$\alpha \frac{\partial^2 T}{\partial x^2} = \frac{\partial T}{\partial t} \quad (0 < x < \delta, t > 0)$$

初始和边界条件

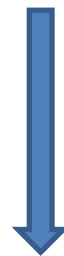
$$T(x, 0) = T_0 \quad (0 \leq x \leq \delta)$$

$$\left. \frac{\partial T(x, t)}{\partial x} \right|_{x=0} = 0 \quad \text{symmetric condition}$$

$$-\lambda \frac{\partial T(x, t)}{\partial x} = h[T(\delta, t) - T_{\infty}]$$

引入过余温度

$$\theta = T(x, t) - T_{\infty}$$



一维平板非稳态导热的控制方程

$$\alpha \frac{\partial^2 \theta}{\partial x^2} = \frac{\partial \theta}{\partial t} \quad (0 < x < \delta, t > 0)$$

初始和边界条件

$$\theta(x, 0) = \theta_0 \quad (0 \leq x \leq \delta)$$

$$\left. \frac{\partial \theta(x, t)}{\partial x} \right|_{x=0} = 0 \quad \text{symmetric condition}$$

$$-\lambda \frac{\partial \theta(x, t)}{\partial x} = h \theta(\delta, t) \quad x = \delta$$

运用分离变量，令

$$\theta = \theta_1(t)\theta_2(x)$$

只是时间的函数，只是空间的函数
代入微分方程

$$\theta_2 \frac{\partial \theta_1}{\partial t} = \theta_1 \alpha \frac{\partial^2 \theta_2}{\partial x^2}$$

两边除以 $\alpha \theta_1 \theta_2$

$$\frac{1}{\alpha \theta_1} \frac{\partial \theta_1}{\partial t} = \frac{1}{\theta_2} \frac{\partial^2 \theta_2}{\partial x^2}$$

引入分离变量

$$\frac{P}{\sqrt{\alpha}}$$

求解常微分方程

$$\frac{1}{\alpha \theta_1} \frac{\partial \theta_1}{\partial t} = -\frac{P^2}{\alpha}$$



$$\theta_1 = e^{-P^2 t}$$

$$\frac{1}{\theta_2} \frac{\partial^2 \theta_2}{\partial x^2} = -\frac{P^2}{\alpha}$$



$$\theta_2 = C_1 \left(\cos \frac{P}{\sqrt{\alpha}} x + C_2 \sin \frac{P}{\sqrt{\alpha}} x \right)$$



通解

$$\theta = e^{-P^2 t} \left[C_1 \left(\cos \frac{P}{\sqrt{\alpha}} x + C_2 \sin \frac{P}{\sqrt{\alpha}} x \right) \right]$$



微分（边界条件中有这一项）

$$\frac{\partial \theta}{\partial x} = \frac{P}{\sqrt{\alpha}} e^{-P^2 t} \left[C_2 \cos \frac{P}{\sqrt{\alpha}} x - C_1 \sin \frac{P}{\sqrt{\alpha}} x \right]$$

代入对称边界条件

$$\left. \frac{\partial \theta(x,t)}{\partial x} \right|_{x=0} = 0 \quad \text{symmetric condition}$$



$$C_2 = 0$$

代入对流换热边界条件 $-\lambda \frac{\partial \theta(x,t)}{\partial x} = h\theta(\delta, t) \quad x = \delta$



$$\frac{P}{\sqrt{\alpha}} e^{-P^2 t} (-C_1 \sin \frac{P}{\sqrt{\alpha}} \delta) = -\frac{h}{\lambda} e^{-P^2 t} (C_1 \cos \frac{P}{\sqrt{\alpha}} \delta)$$



令

$$\mu = \frac{P\delta}{\sqrt{\alpha}}$$

$$\tan \mu = \frac{h/\lambda}{P/\sqrt{\alpha}} = \frac{B_i}{\mu}$$
$$c \tan \mu = \frac{\mu}{B_i}$$

← 超越方程

通解的形式

$$\begin{aligned}\theta &= e^{-P^2 t} \cos \frac{P}{\sqrt{\alpha}} x \\&= \sum_{n=1}^{\infty} C_n e^{-P^2 t} \cos \frac{P}{\sqrt{\alpha}} x \\&= \sum_{n=1}^{\infty} C_n e^{-\frac{P^2 \delta^2}{\alpha} \frac{\alpha t}{\delta^2}} \cos \frac{P \delta}{\sqrt{\alpha}} \frac{x}{\delta} \\&= \sum_{n=1}^{\infty} C_n e^{-\mu^2 F_o} \cos(\mu \eta)\end{aligned}$$

根据正交函数
性质及傅里叶
级数展开



$$C_n = \frac{2 \sin \mu_n}{\mu_n + \cos \mu_n \sin \mu_n}$$

温度与无量纲参数的关系

$$\theta = \frac{T - T_{\infty}}{T_0 - T_{\infty}} = \text{func}(Fo(t), \frac{x}{\delta}, Bi)$$

2014/03/18

内容提要

◆半无限平板的非稳态导热

◆对流传热理论基础

◆教科书5.1

1. 影响对流传热的因素

2. 表面传热系数的确定

◆教科书5.3 对流边界层

1. 速度边界层

2. 热边界层

3. 层流和湍流

4. 雷诺数

5. 努赛尔数

半无限平板的温度变化

$$\alpha \frac{\partial^2 T}{\partial x^2} = \frac{\partial T}{\partial t}$$



$$\eta \equiv \frac{x}{\sqrt{4\alpha t}}$$

$$\frac{d^2 T}{d\eta^2} = -2\eta \frac{dT}{d\eta}$$



$$\frac{T - T_s}{T_0 - T_s} = \frac{2}{\sqrt{\pi}} \int_0^\eta \exp(-\eta^2) d\eta \equiv \operatorname{erf}(\eta)$$

$$\begin{aligned}
 q_x &= -\lambda \frac{\partial T}{\partial x} = -\lambda(T_0 - T_s) \frac{\partial \operatorname{erf}(-\eta^2)}{\partial x} \\
 &= -\lambda(T_0 - T_s) \frac{2}{\sqrt{\pi}} \frac{\partial \int_0^\eta \exp(-\eta^2) d\eta}{\partial \eta} \frac{\partial \eta}{\partial x} \\
 &= -\lambda(T_0 - T_s) \frac{2}{\sqrt{\pi}} \exp(-\eta^2) \frac{1}{2\sqrt{\alpha t}} \\
 &= \lambda \frac{T_s - T_0}{\sqrt{\pi \alpha t}} \exp(-\eta^2) \\
 &= \lambda \frac{T_s - T_0}{\sqrt{\pi \alpha t}} \exp\left[-\frac{x^2}{4\alpha t}\right]
 \end{aligned}$$

教科书 P135 Eq 3-42

通过表面的热流密度

$$q_{x=0} = \lambda \frac{T_s - T_0}{\sqrt{\pi \alpha t}}$$

一段时间内表面上的传热量

$$\begin{aligned} Q &= A \int_0^t \lambda \frac{T_s - T_0}{\sqrt{\pi \alpha t}} dt \\ &= A \frac{\lambda (T_s - T_0)}{\sqrt{\pi \alpha}} 2\sqrt{t} = A \sqrt{\frac{t}{\pi}} \sqrt{\rho c \lambda} (T_s - T_0) \end{aligned}$$

表面传热系数

$$h = - \frac{\lambda}{\Delta T_s} \left. \frac{\partial T(x, t)}{\partial x} \right|_{x=0}$$

◆表面传热系数，由流体导热系数，贴壁处壁面法线方向上的流体温度变化率，壁面与流体的温度差决定

◆速度边界层的主要表现形式为**表面摩擦**，影响表面摩擦系数

◆热边界层的主要表现形式为**对流换热**，会影响对流换热系数

对流传热在工业中的应用

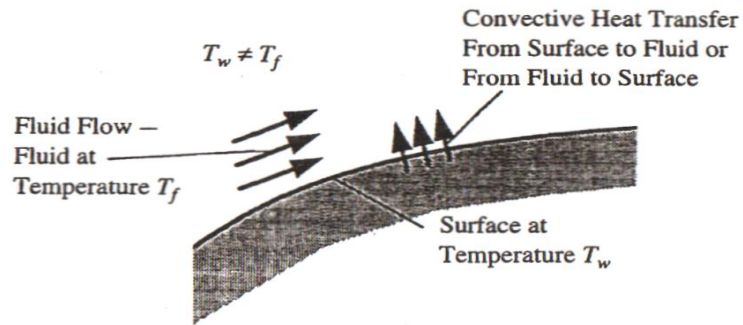


FIGURE 1.1
Convective heat transfer.

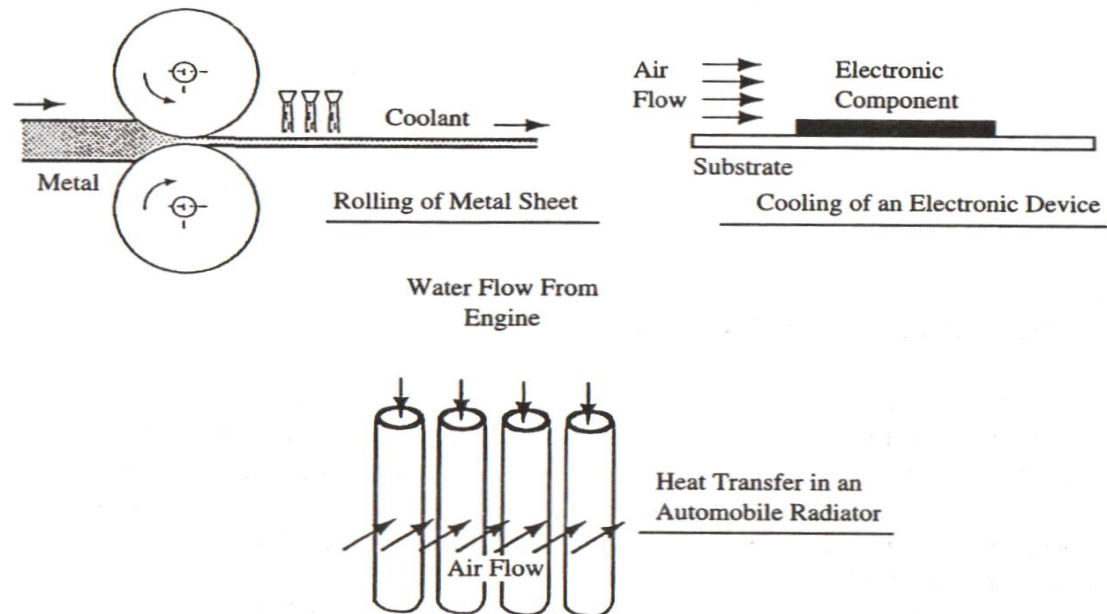
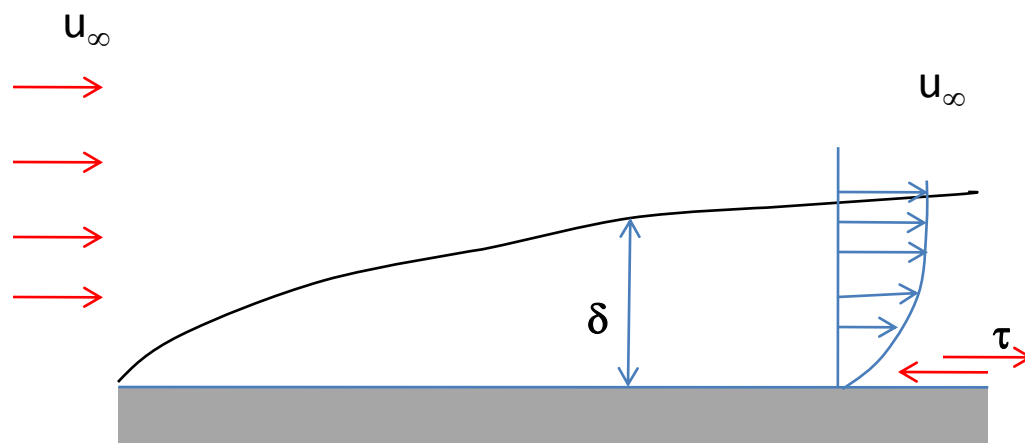
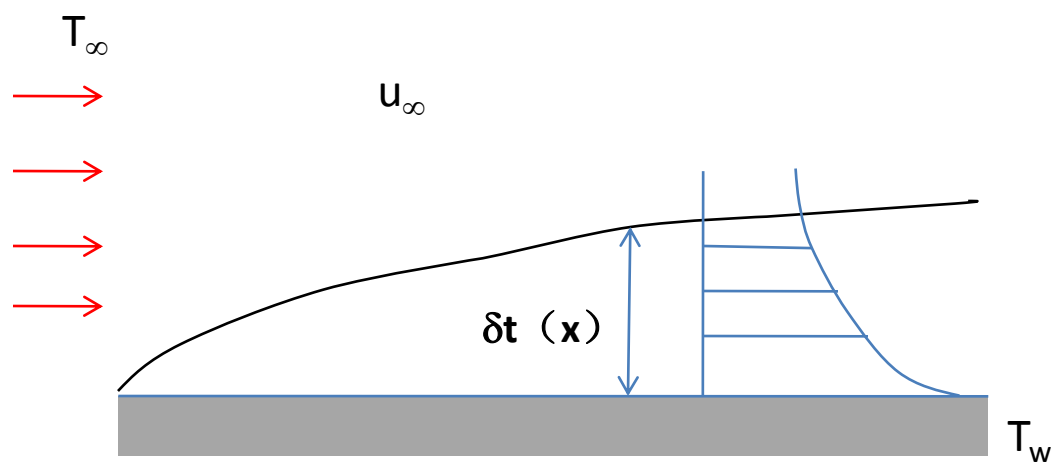


FIGURE 1.2
Some situations that involve convective heat transfer.

速度边界层



热边界层



层流到湍流的过渡

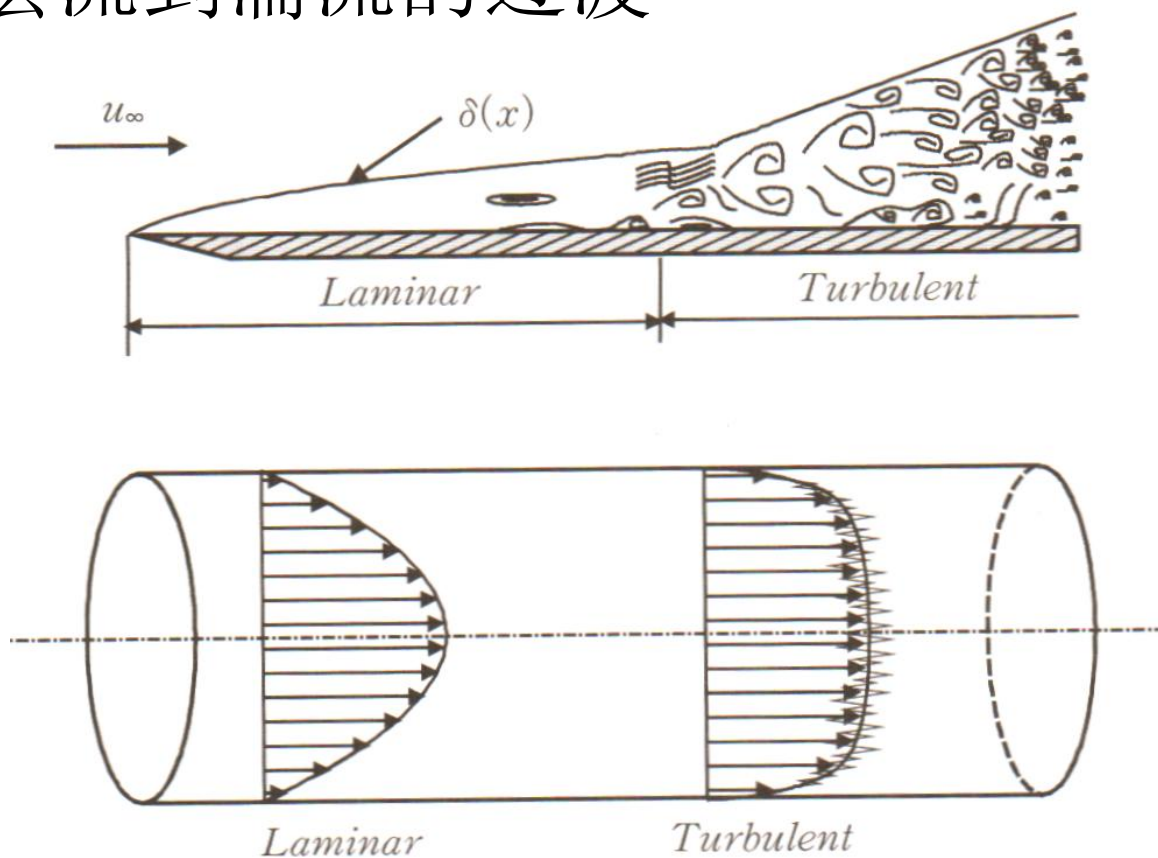


図 4.2 層流から乱流への遷移

努赛尔数Nu的物理意义

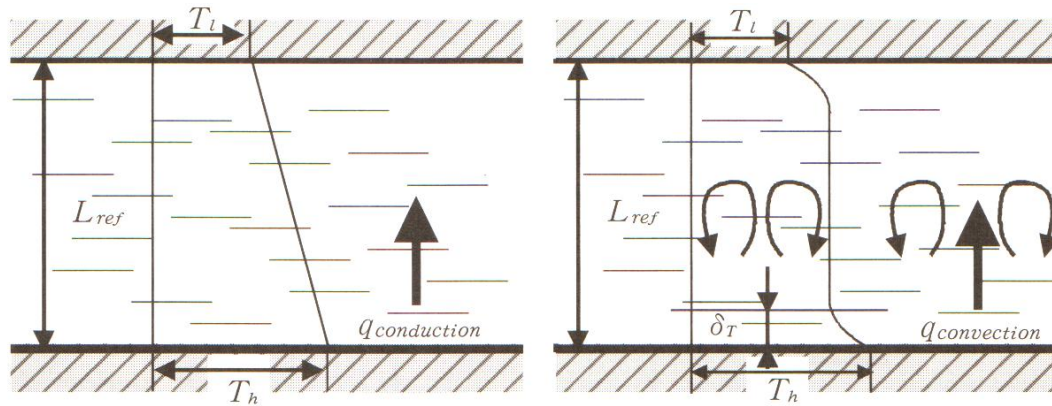


図 4.3 ヌッセルト数の意味

由导热引起的流体热流密度

$$\frac{q_{convection}}{q_{conduction}} = \frac{h(T_h - T_l)}{k_f \frac{(T_h - T_l)}{L_{ref}}} = \frac{hL_{ref}}{k_f} = Nu_{ref}$$

$$q_{conduction} = k_f \frac{(T_h - T_l)}{L_{ref}}$$



$$q_{convection} = h(T_h - T_l)$$