I. Backbone exponents.

 \star backbones defined as current-carrying bonds (sites) in bus-bar geometry. $N_b \sim L^{d-X_b}$.

 N_b — number of backbones, L — finite linear system size; d — spatial dimensionality; X_b — backbone scaling dimension.

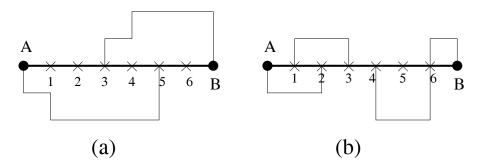


Burning Monte Carlo algorithm.

 \star path crossing exponents X_k . probability P_k of two points connected by at least k independent paths: $P_k \sim L^{-2X_k}$.

$$\Downarrow X_b = X_2$$

current Monte Carlo algorithm.



† † much more efficient than burning procedures.

* results.

1): 2D q-state Potts model (including Baxter-Wu model). bond probability p set at random cluster fixed point p ($p = 1 - e^{-K}$), and at Potts cluster fixed point p (p = 1).

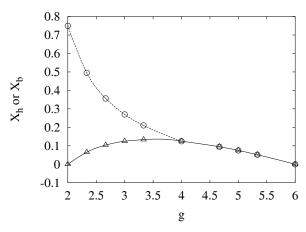
| q | g | Mod. | $X_h^{(r)}$ | $X_b^{(r)}$ | $X_h^{(p)}$ | $X_b^{(p)}$ |
|---|------|---------|-------------|-------------|-------------|-------------|
| 0 | 2 | q = 0 P | 0 | 3/4? | -3/16 | |
| 1 | 8/3 | Perc. | 5/48 | 0.3566(2) | 0 | 0 |
| 2 | 3 | Ising | 1/8 | 0.2696(3) | 5/96 | 0.0520(3) |
| 3 | 10/3 | q = 3 P | 2/15 | 0.2105(3) | 7/80 | 0.0871(7) |
| 4 | 4 | q=4 P | 1/8 | 0.126(1) | 1/8 | 0.1246(5) |
| 4 | 4 | B.W. | 1/8 | 0.136(7) | 1/8 | 0.1239(8) |
| 2 | 14/3 | Trc. I. | 3/40 | 0.0760(15) | 0.0752(3) | 0.0753(8) |

Note: 1), q=4 Potts model is at tricritical point (K=1.45790(1), D=2.478438(2)); 2), $X_b=2/3$ for q=0 Potts obtained from Eden tree.

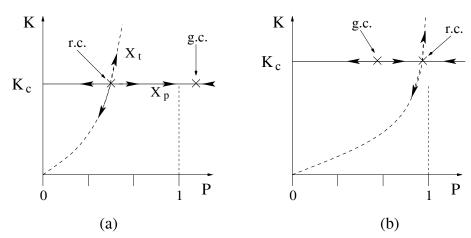
* Q1: can one exactly calculate X_b in 2D?

Fit formula:

$$P = L^{-2X}(a + bL^{y_1} + cL^{-2} + dL^{-3}) \quad \text{or} \\ P = L^{-2X}(a + b/\ln L + c + b/\ln^2 L + dL^{-2}) \text{ (P.4 and BW)}.$$



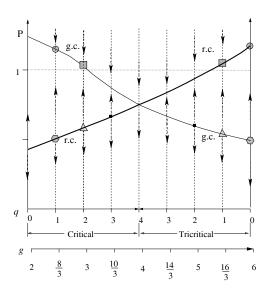
- $\dagger \dagger X_b = X_h$ for $g \ge 4$, $X_b > X_h$ for g < 4.
- 2): 3D Ising $X_b = 0.829(4)$. 3D percolation and tricritical Ising $X_b = ?$
- II. 'Geometric cluster' fixed points.
 - † † $X_b = X_h$? depends on 'red-bond' exponent $X_r > 2$?;
 - † † Scaling argument $\Rightarrow X_r = X_p$ (RG exponent along bond-probability direction.
- \star 2D q-state Potts model.
 - e.g., Ising model (a) and tricritical Blume-Capel model (b) on square lattice.



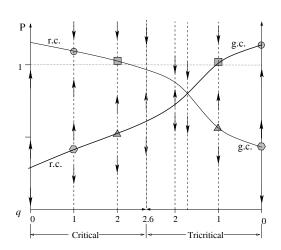
† † at r.c., $X_p > 0$ for critical branch, $X_p > 0$ for tricritical branch, and $X_p = 0$ for 4-state Potts model.

• From Kac formula and other assumptions, we obtain following tables (last two pages).

RG flows in p-q space:

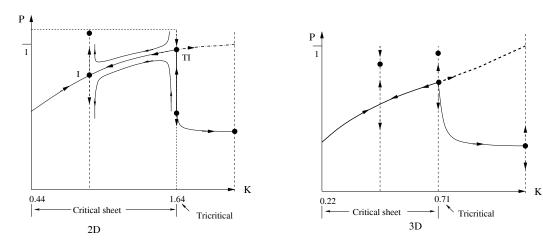


 \star for 3D Ising model $X_p=0.767(2),~X_h^{(g)}=0.14(1)$?; for tricritical Blume-Capel model $X_P>0$, so that RG flows:



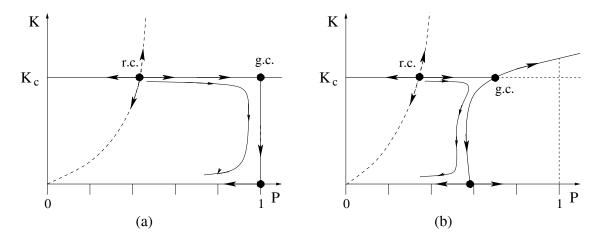
Q2, for 3D q=0 Potts, are followings correct? at criticality, $X_t=3$ and $X_h=0$? at tricriticality, $X_t=0.5183$ and $X_h=0.14(1)$?

Q3, for Blume-Capel model, RG flows and percolation threshold are show in following figures, are they correct?



Q4, for tricritical 3D q=2 Potts, since it behaves mean-field-like, can one calculate X_b ?

Q5, concerning Ising models on triangular and simple-cubic lattices, (a) and (b), respectively, are following RG flows correct?



† † for Ising clusters (P = 1), percolation thresholds coincide with K_c for triangular lattice, but not for simple-cubic lattice.